

Havse 2011

Hamburg



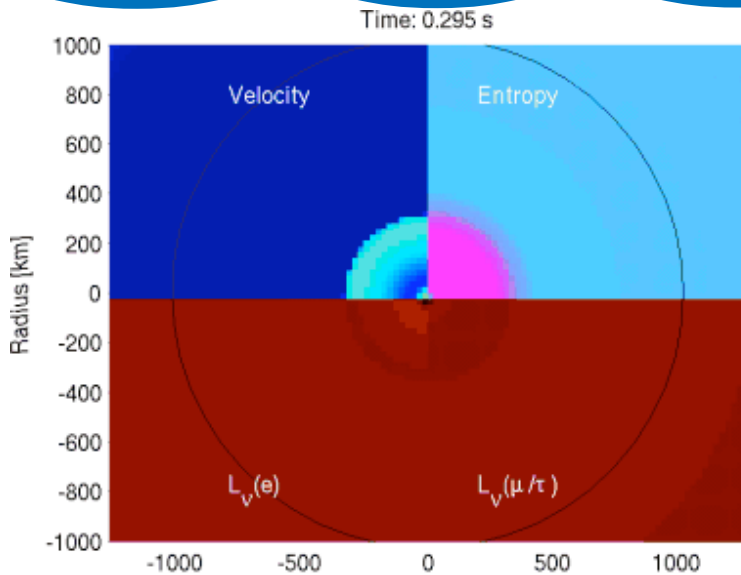
Microphysics of Supernova Core

M. Liebendörfer
University of Basel

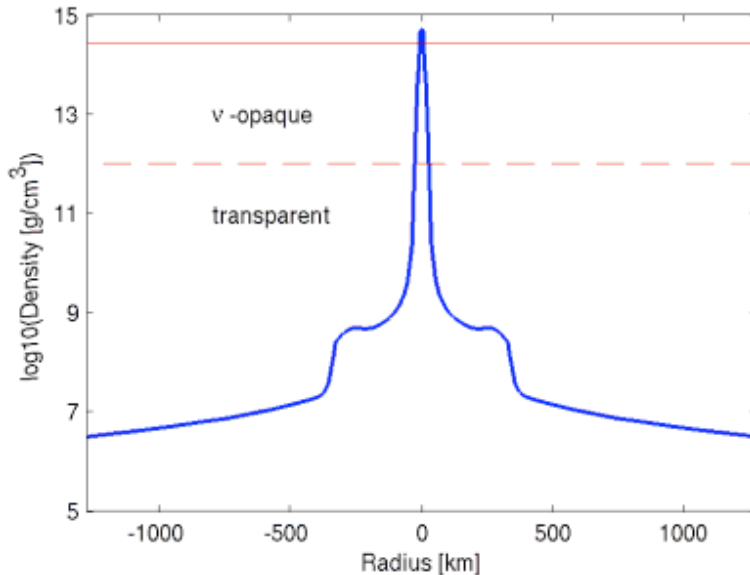
- Matter and neutrinos in the collapse phase
- Matter and neutrinos in the postbounce phase
- New 3D models of 11 Ms, 15 Ms and 40 Ms progenitor

R. Cabezón
M. Hempel
R. Käppeli
A. Perego
F.-K. Thielemann
N. Vassetz
S. C. Whitehouse

The Explosion Mechanism(s)



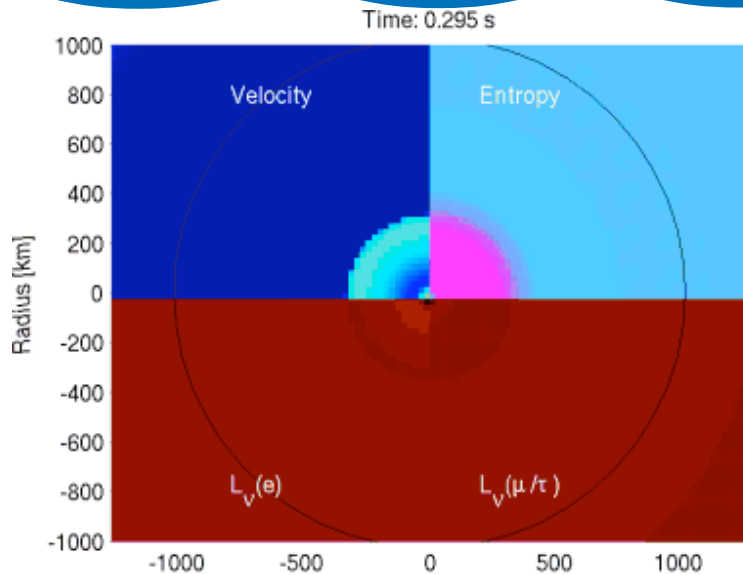
Keep in mind:
The explosion is
only a surface
effect on the
proton neutron
star!



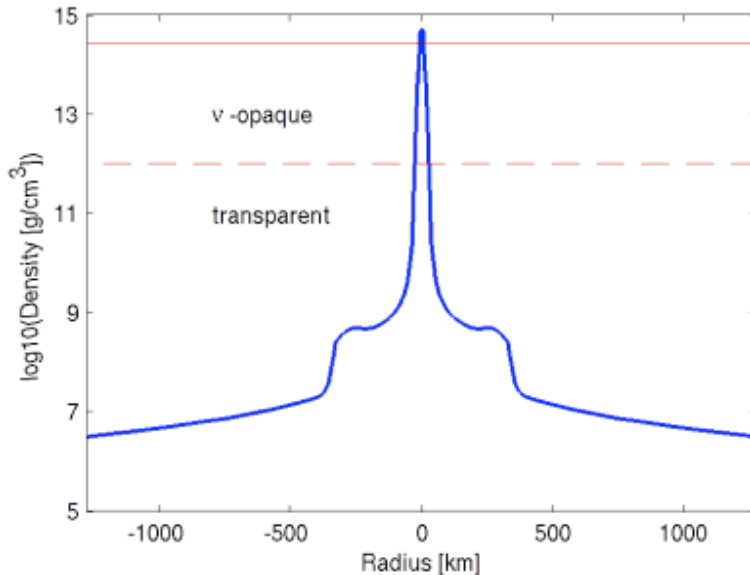
Different
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mechanisms
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- prompt hydrodynamic
- delayed, ν -driven
- magneto-rotational
- acoustic
- magnetoviscous/sonic
- phase transition in NS

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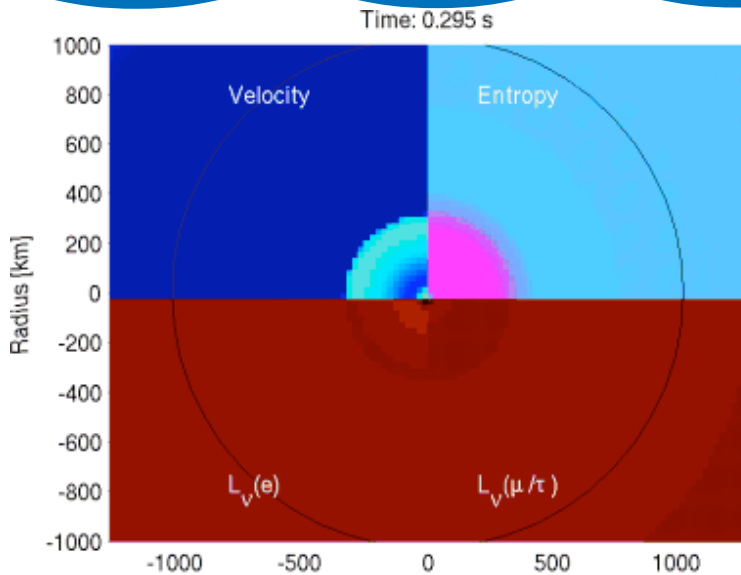
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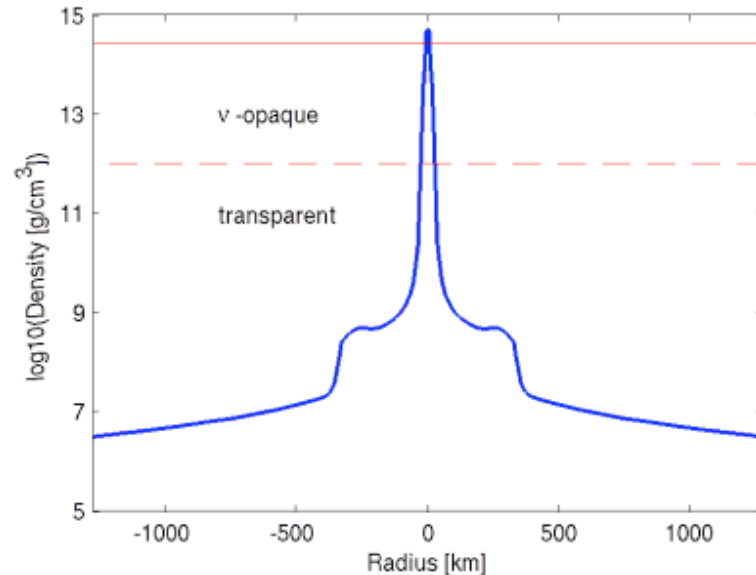
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- prompt hydrodynamic (e.g. Baron et al. 1985)
- delayed, ν -driven (Colgate 1966, Bond 1979, ... Marek & Janka 2009)
- magneto-rotational (Bisnovatyi-Kogan 1976, Leblanc & Wilson 1979, ...)
- acoustic (Burrows et al. 2006)
- magnetoviscous/sonic (Akiyama et al. 2003, Thompson et al. 2003, Socrates et al. 2005)
- phase transition in NS (Migdal et al. 1971, ... Sagert et al. 2009)

The Explosion Mechanism(s)



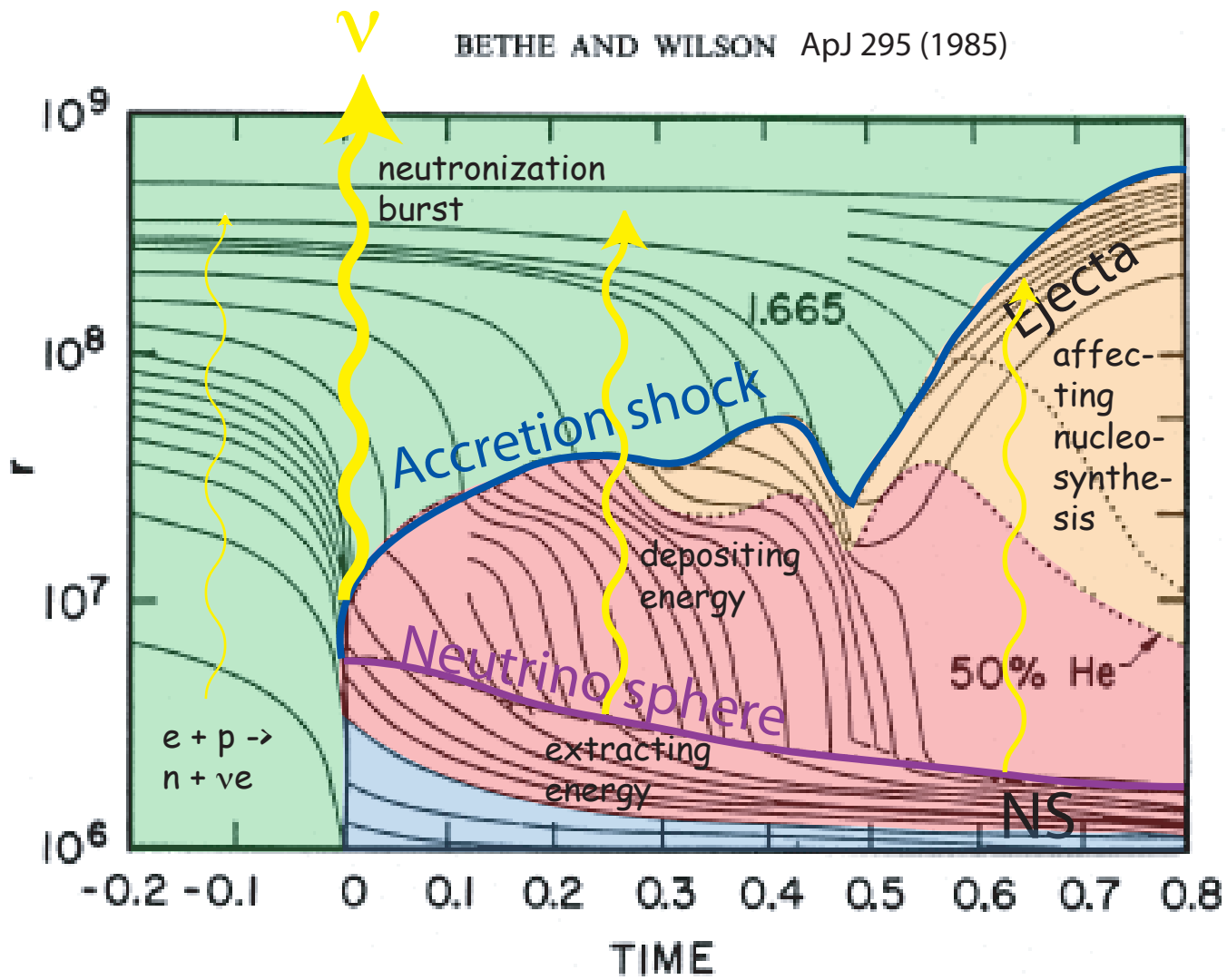
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Sagert et al. 2009)

Matter conditions



Ensemble
of nuclei

Cool bulk
nuclear matter

Hot dissociated
matter

Freeze-out
of nuclei

collapse phase || postbounce accretion phase | explosion phase
bounce

Neutrino-matter interactions

Bruenn (1985)
Raffelt (2001)



Description:

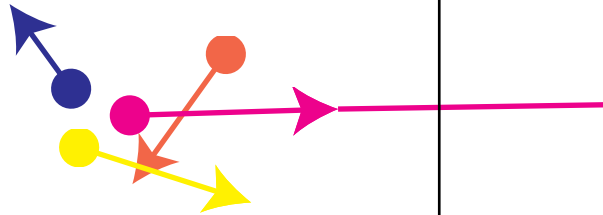
Number Sphere

Energy Sphere

Transport Sphere

Emission &
Absorption

High matter density

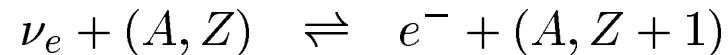
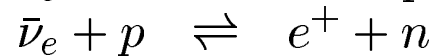
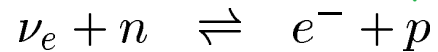


Can be Pauli-
blocked in
diffusive
regime

Low matter density

Production/Annihilation

Electron or Positron capture

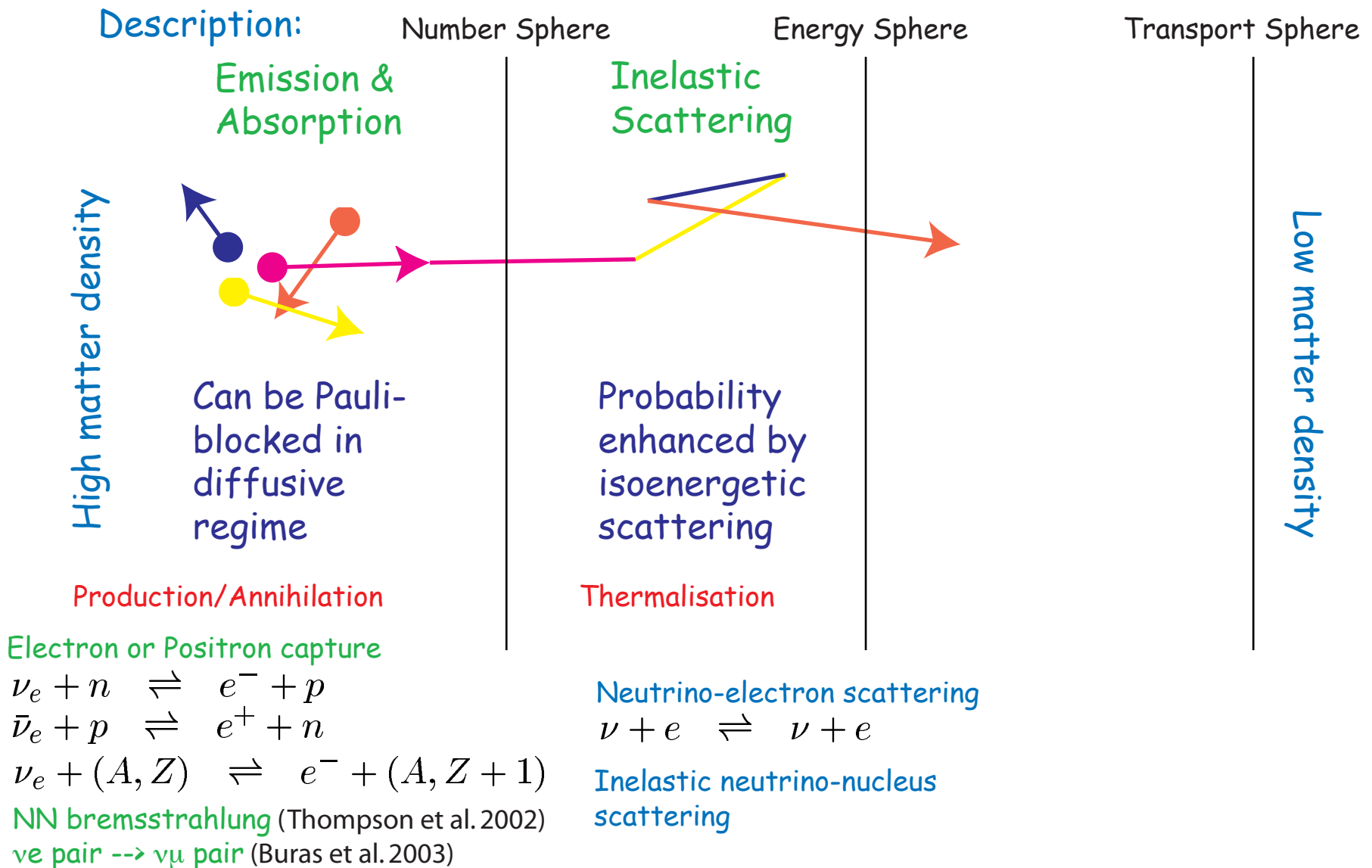


NN bremsstrahlung (Thompson et al. 2002)

ν_e pair \rightarrow ν_μ pair (Buras et al. 2003)

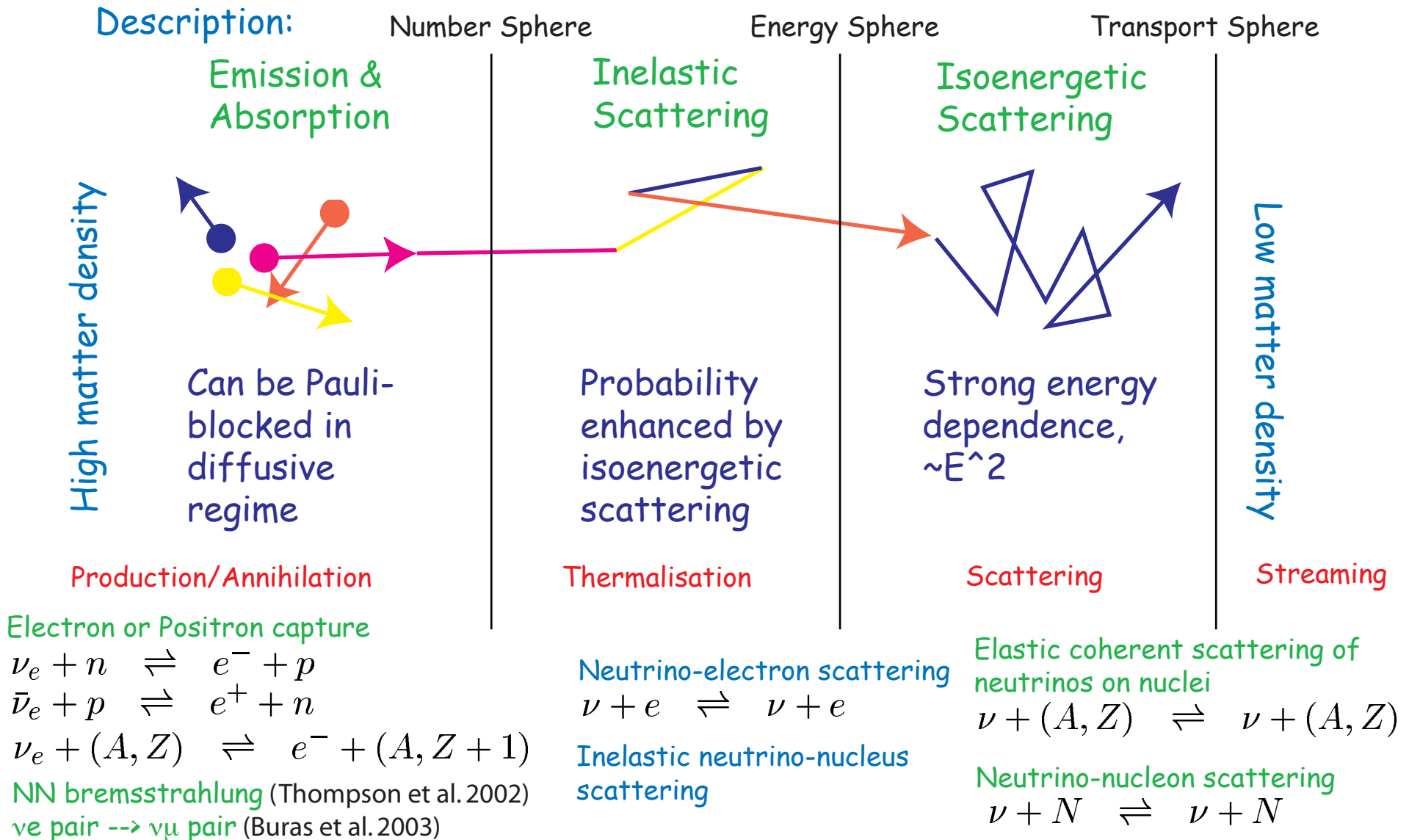
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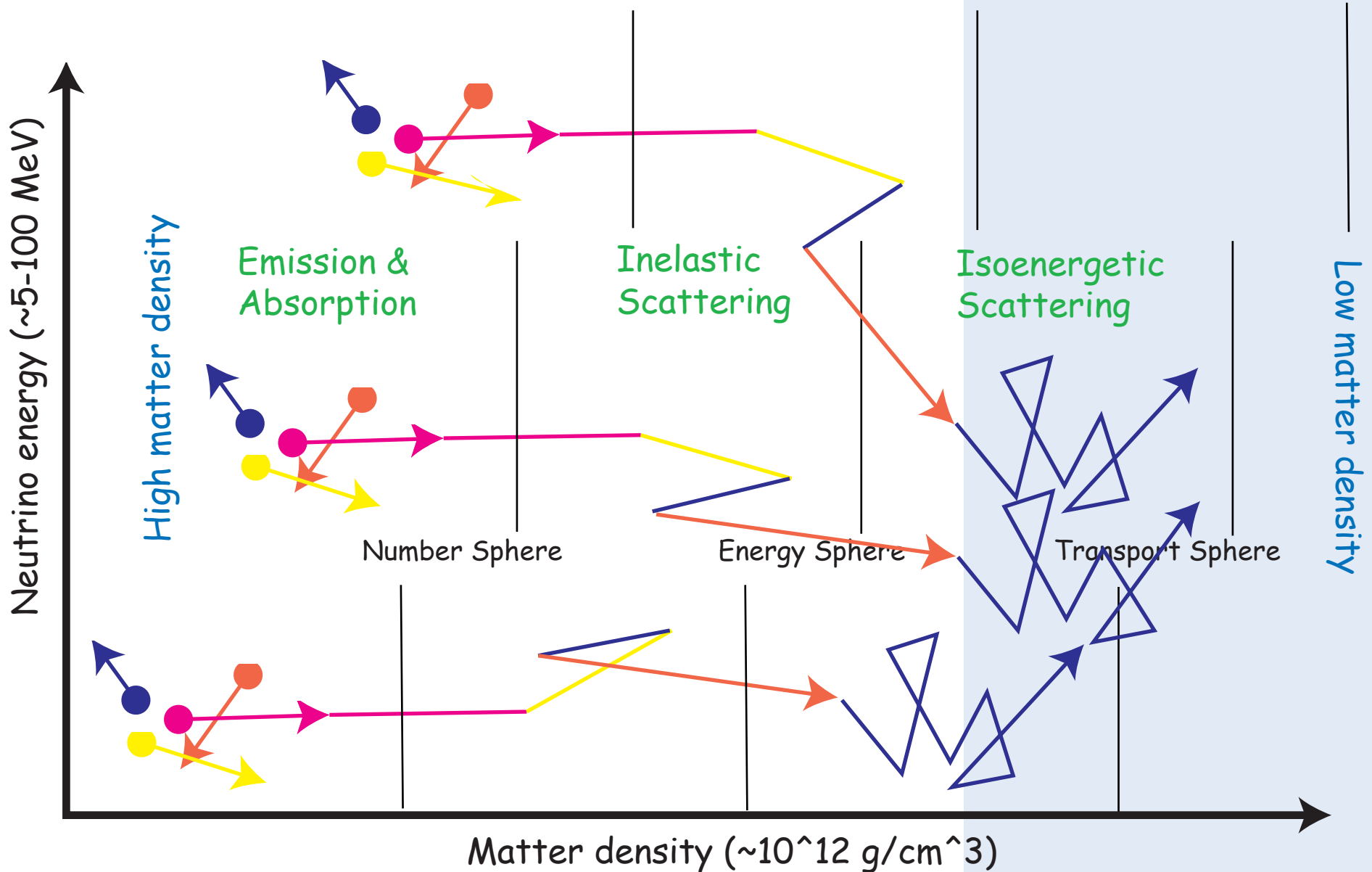


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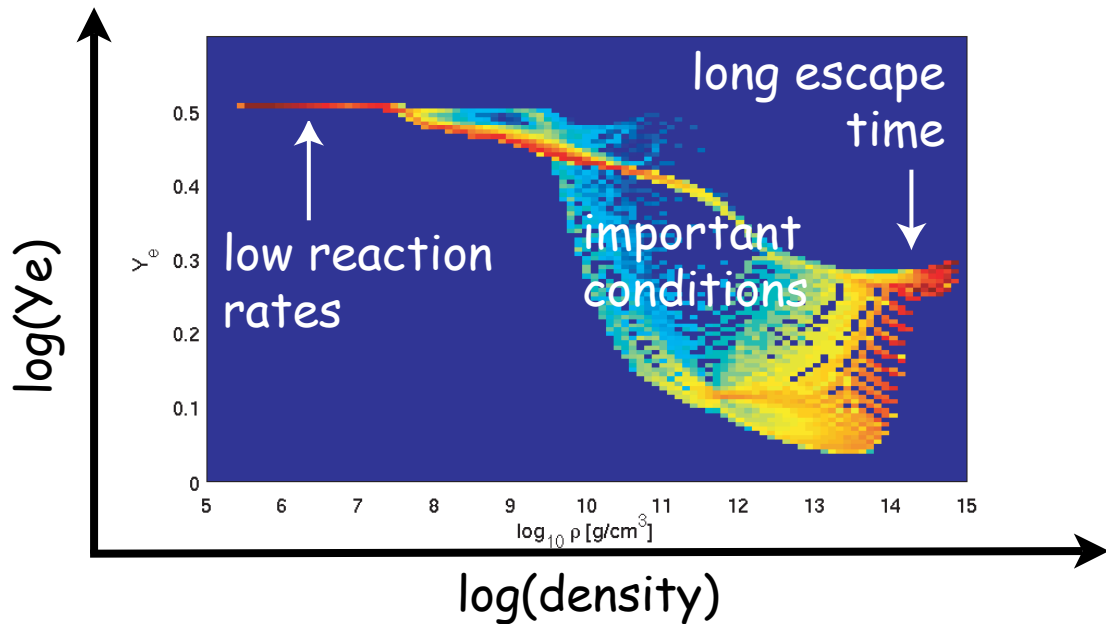
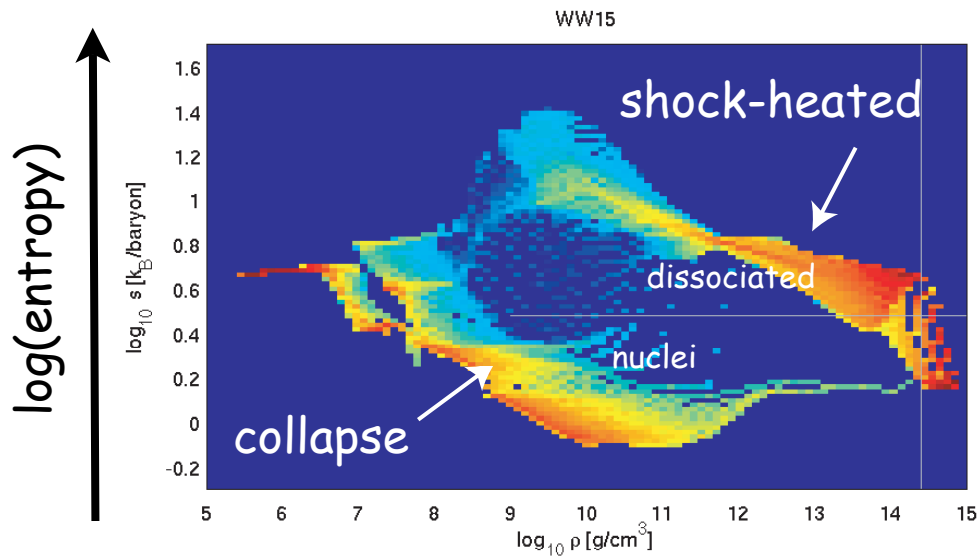
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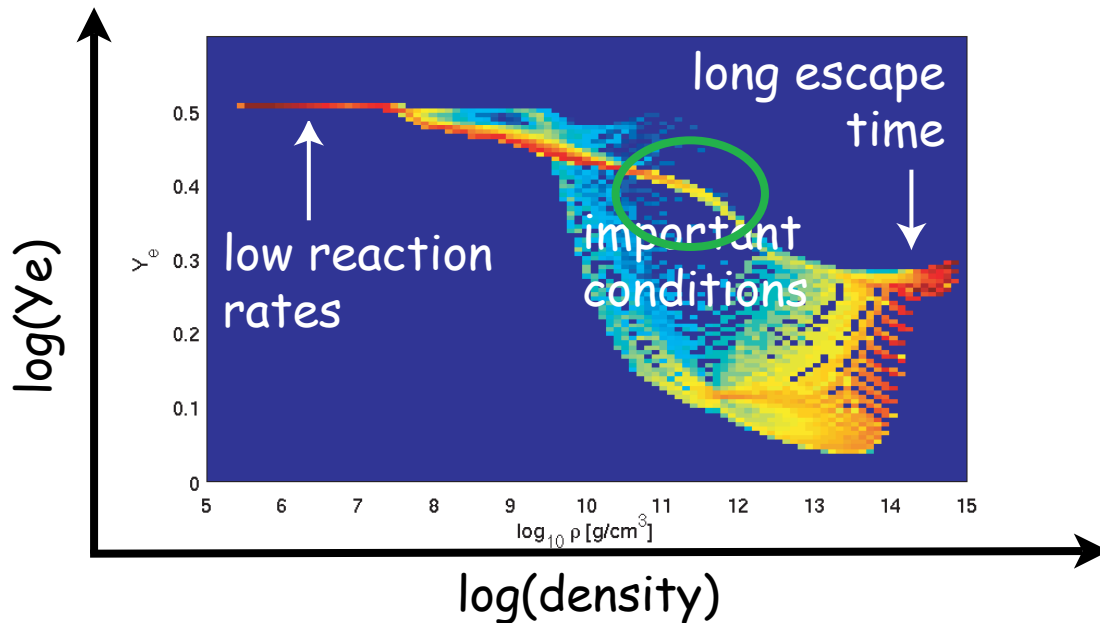
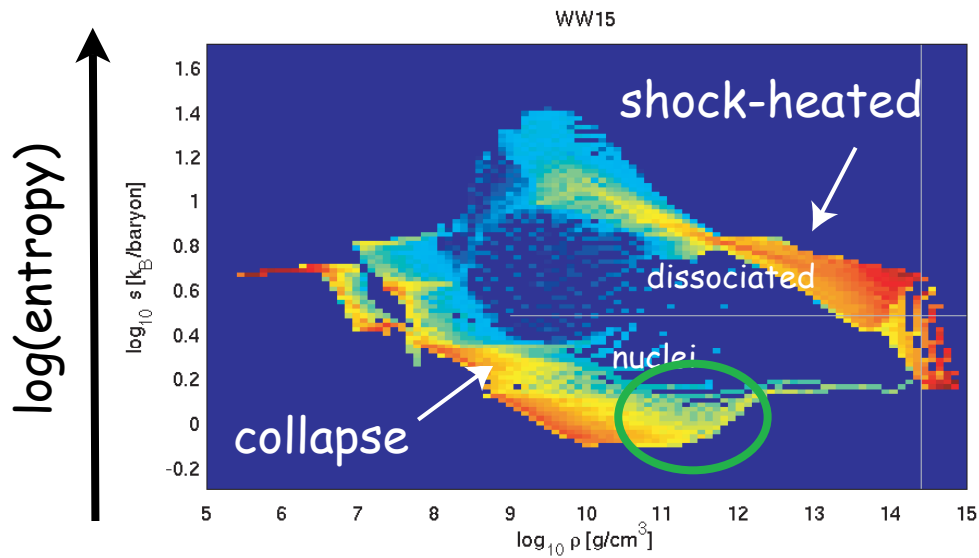
Energy-dependent neutrino transport



Relevant ν -matter interactions



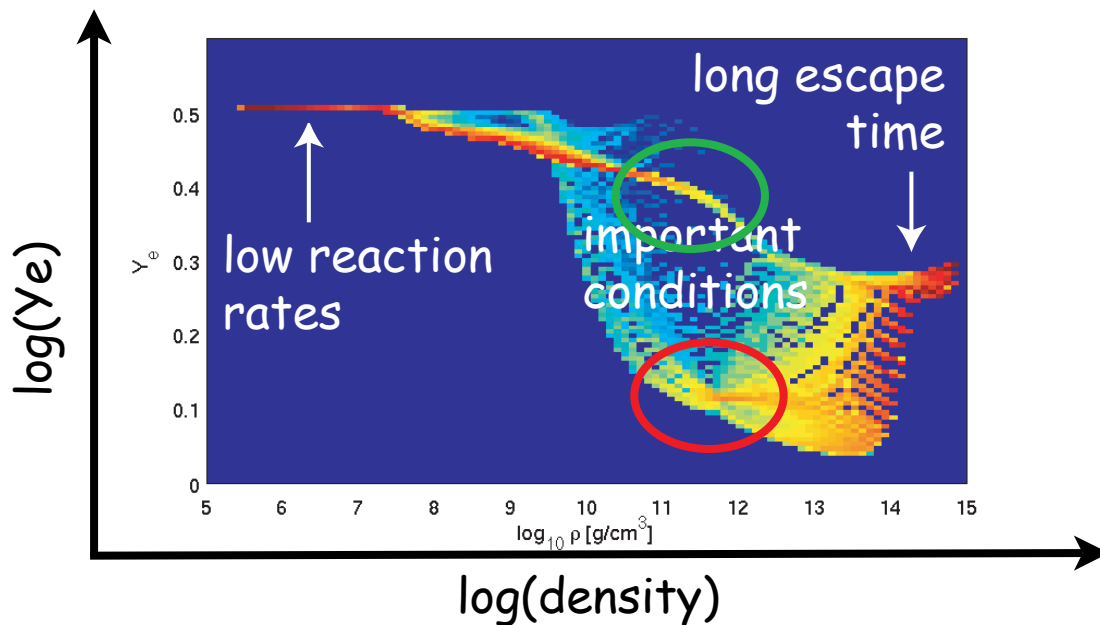
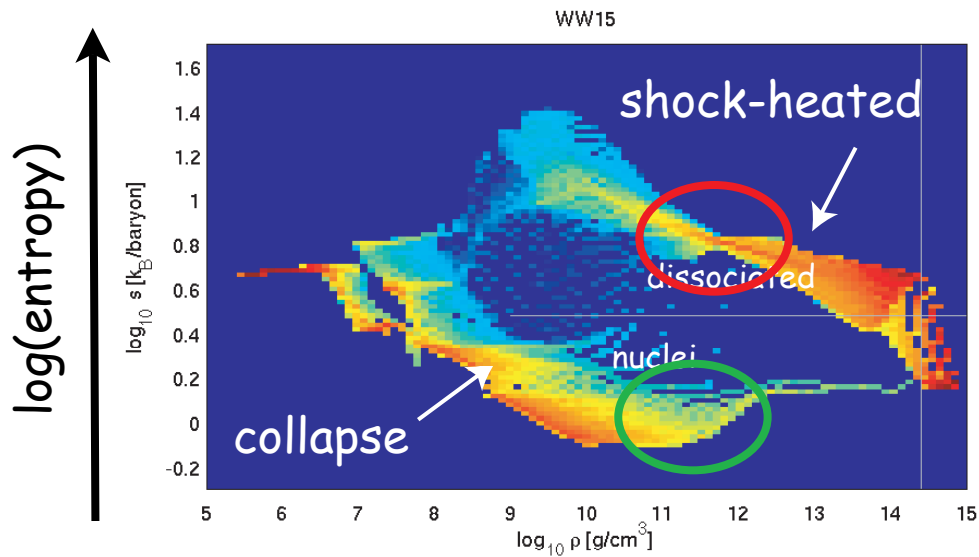
Relevant ν -matter interactions



The conditions around the neutrino spheres are marked in

green ...collapse

Relevant ν -matter interactions



The conditions around the neutrino spheres are marked in

green ... collapse

red ... postbounce

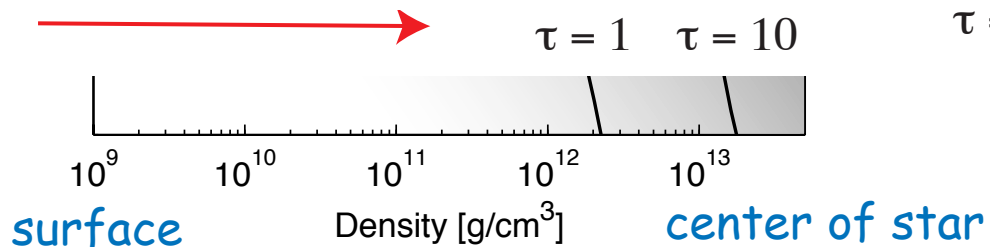
Physics for the collapse phase

Bethe (1990)
mean free path:

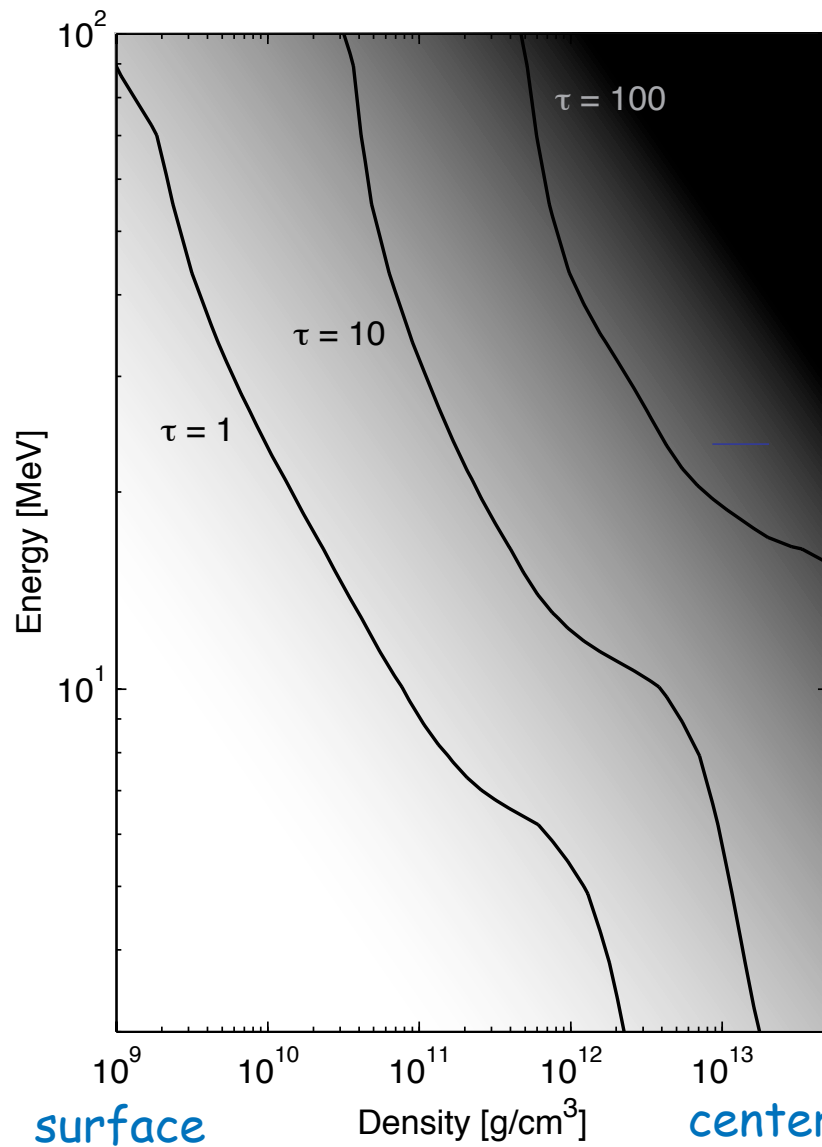
$$\lambda_\nu = 1.0 \times 10^8 \rho_{12}^{-1} [(N^2/6A)X_h + X_n]^{-1} \epsilon_\nu^{-2} \text{ cm} .$$

Optical depth:

$$\tau = \int dr/\lambda$$



Deleptonization



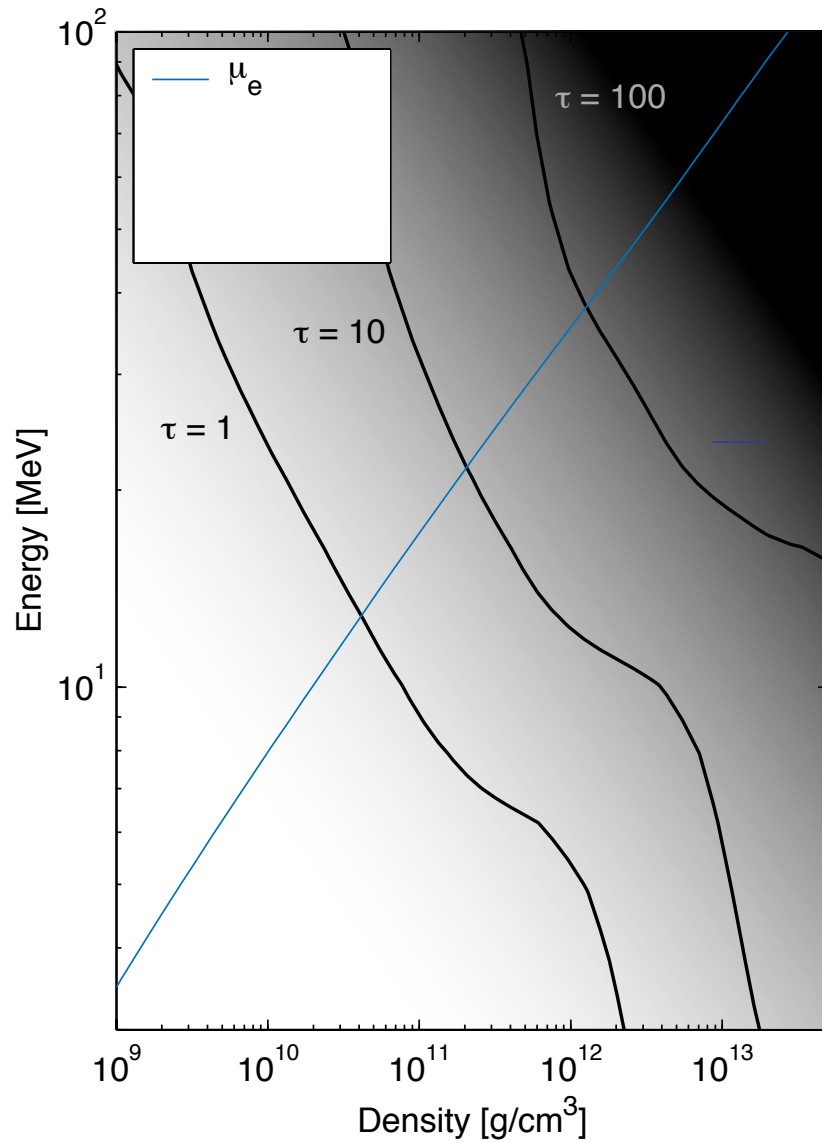
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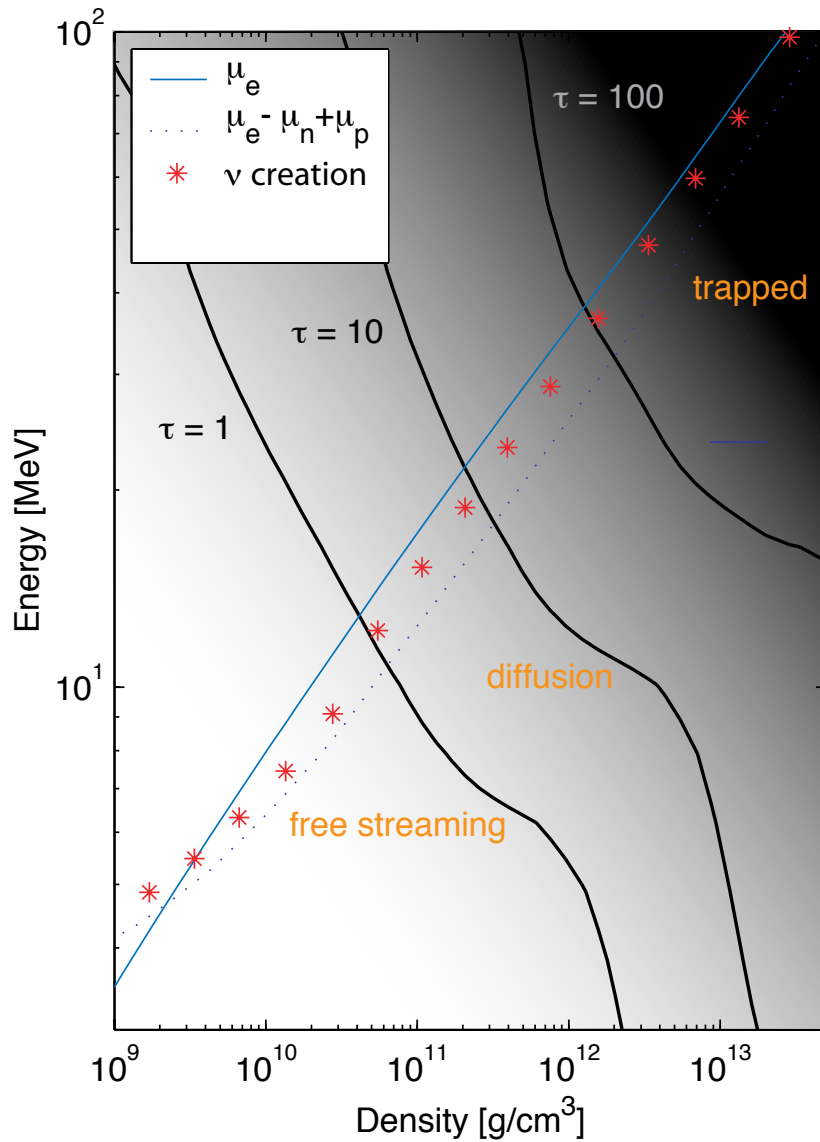
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Deleptonization



(Martinez-Pinedo, Liebendoerfer, Frekers, 2006)

Deleptonization

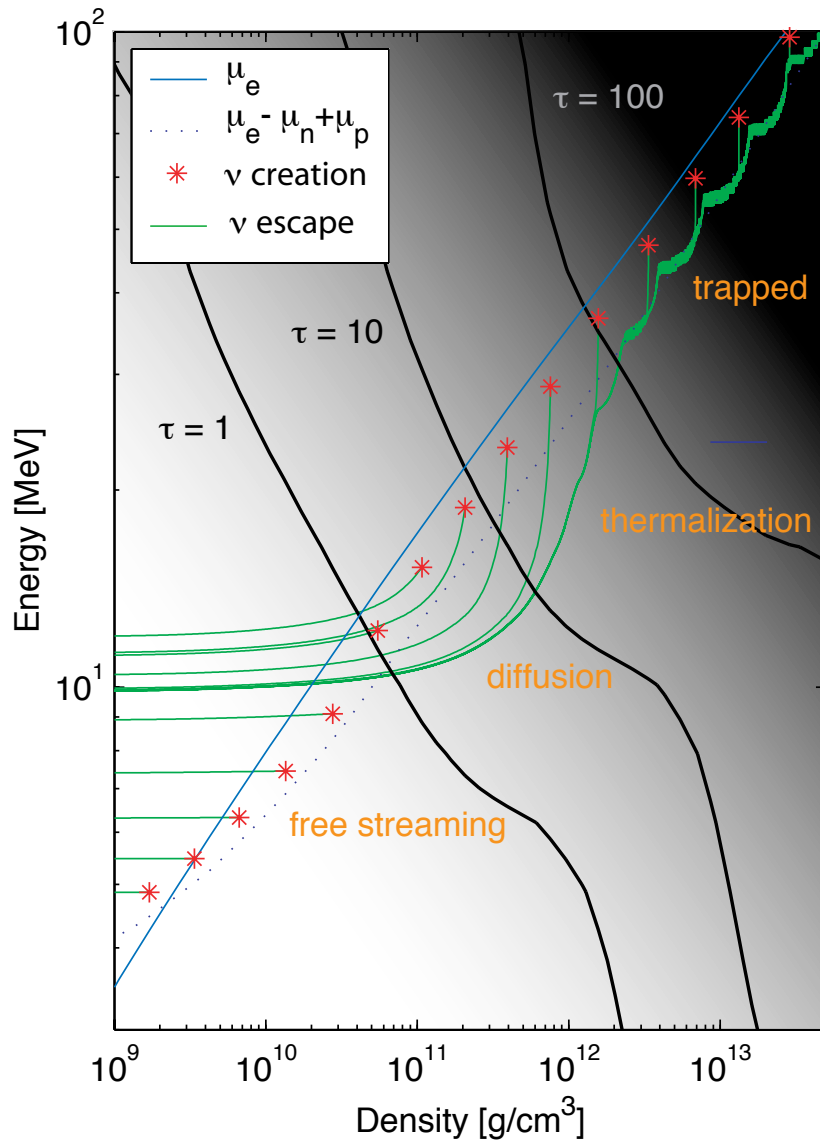


$$\mu_\nu \sim \mu_e - \mu_n + \mu_p$$

$$\frac{\Delta s}{\Delta t} = - \frac{\Delta Y_e}{\Delta t} \frac{\mu_e - \mu_n + \mu_p - E_\nu^{esc}}{T}$$

(Martinez-Pinedo, Liebendoerfer, Frekers, 2006)

Deleptonization



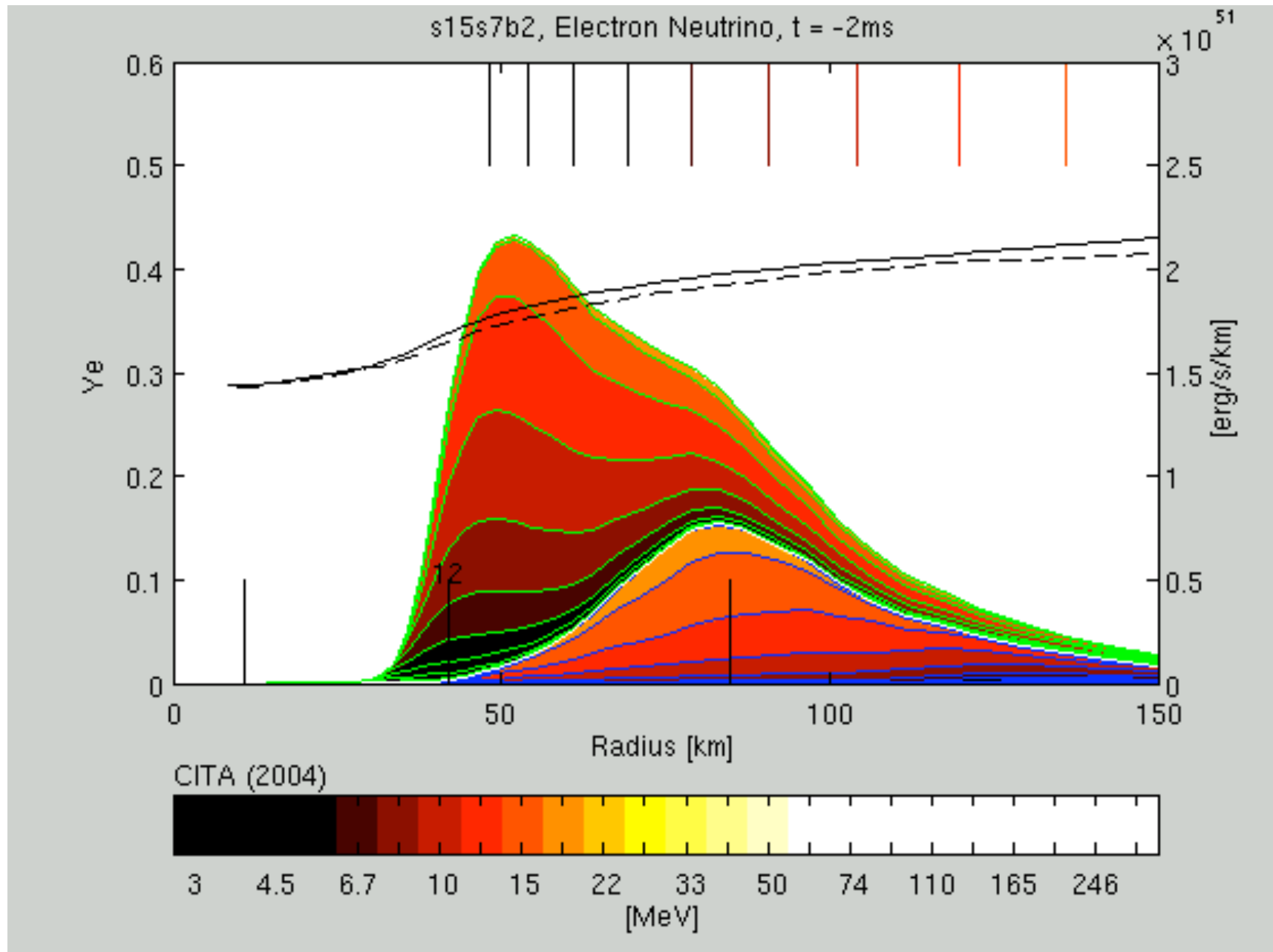
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$$\frac{\Delta s}{\Delta t} = - \frac{\Delta Y_e \mu_e - \mu_n + \mu_p - E_\nu^{esc}}{\Delta t T}$$

ν 's escape directly:
→ entropy decrease

ν 's thermalise:
→ entropy increase

Histogram of electron neutrino emission

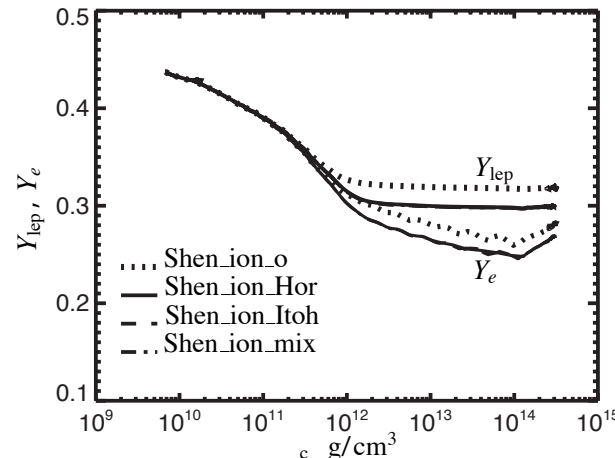
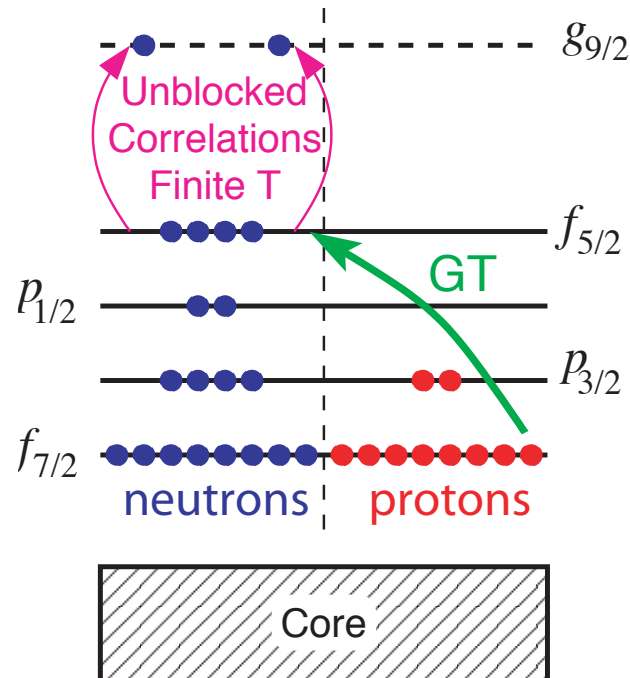


Physics for the collapse phase

Electron capture rates
 - ensemble of nuclei
 with individual rates

Thermalisation
 - neutrino-electron
 scattering, inelastic
 scattering with nuclei

Opacity
 - coherent scattering
 - ion-ion correlations
 - clustering in phase
 transition
 - unresolved: how to
 scatter on ensemble of
 nuclei?



Langanke et al. 2003, Hix et al.
 2003, Martinez-Pinedo,
 Liebendörfer, Frekers 2006

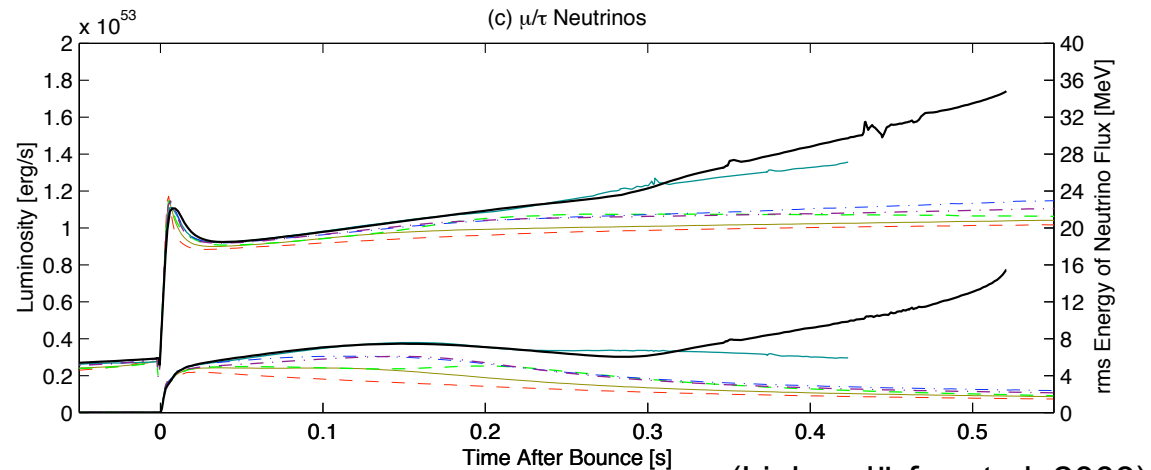
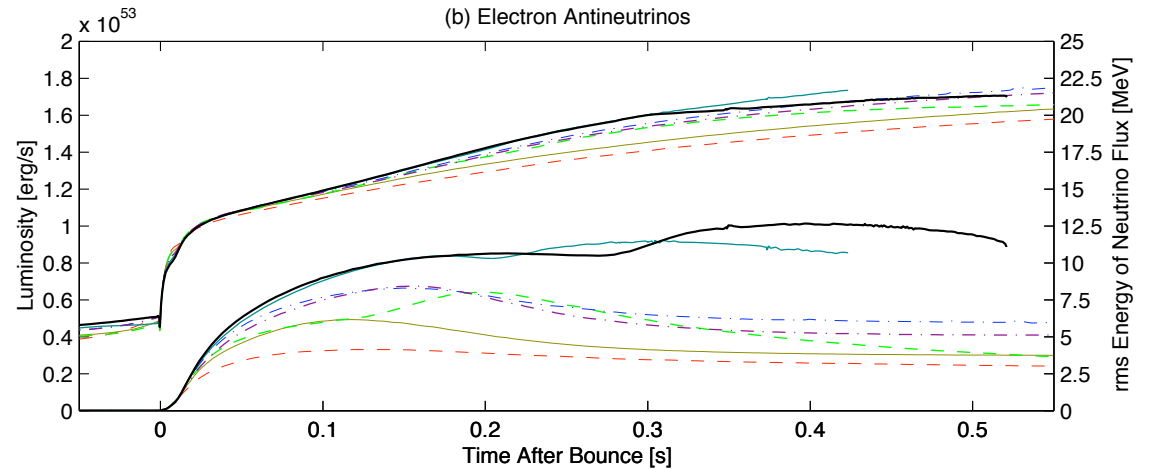
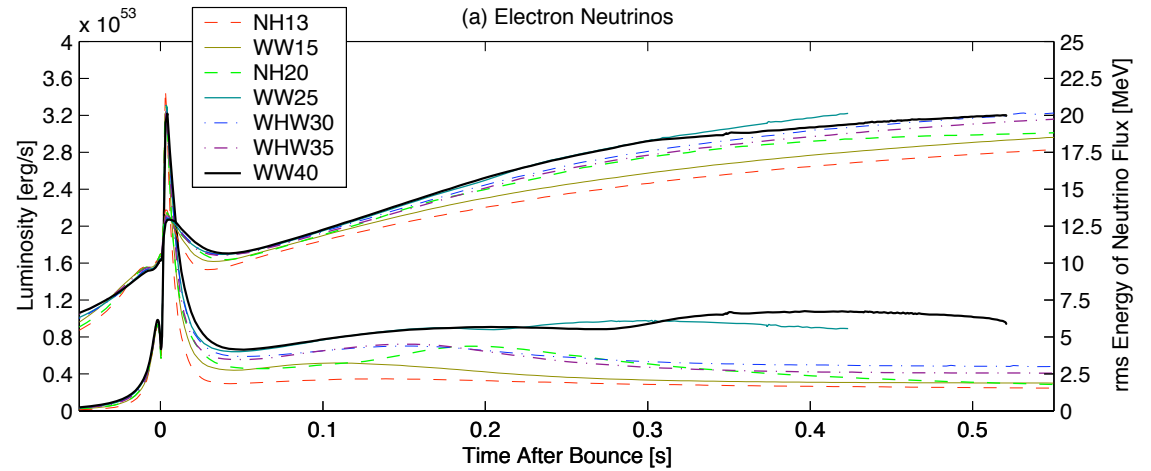
Myra & Bludman, Bruenn 1989,
 Langanke et al. 2008

Itoh 1975, Horowitz 1997, Bruenn
 & Mezzacappa 1997, Watanabe
 2004, Horowitz et al. 2004,
 Botvian & Mishustin 2005, Marek
 et al. 2005

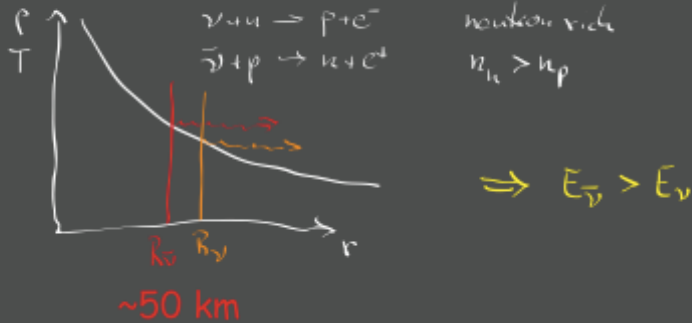
Structure and
 dynamics of nuclei
 are key!

Neutrinos from the postbounce phase

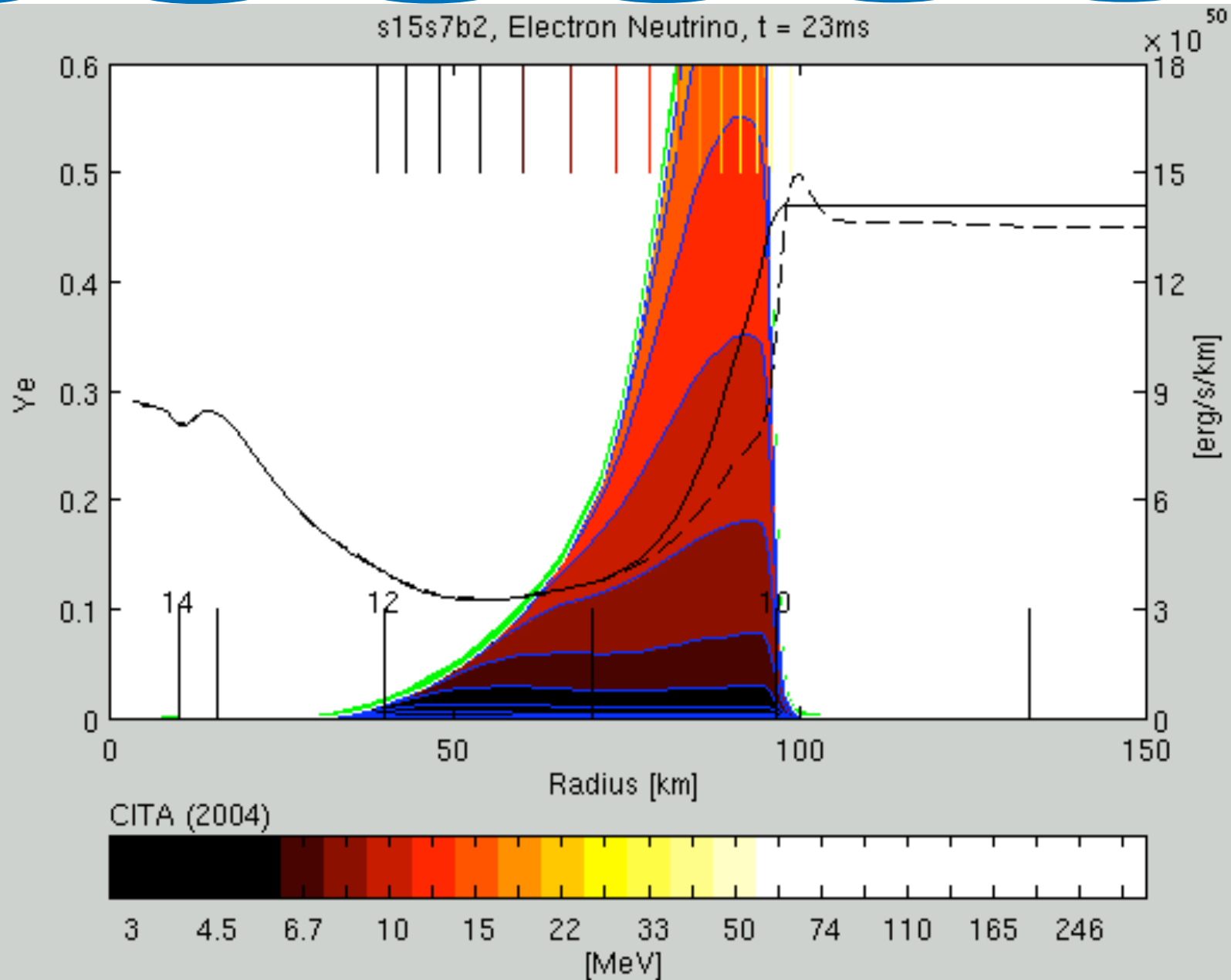
The neutrino luminosities reflect the accretion rate and the thermodynamic conditions at the **neutrinospheres**



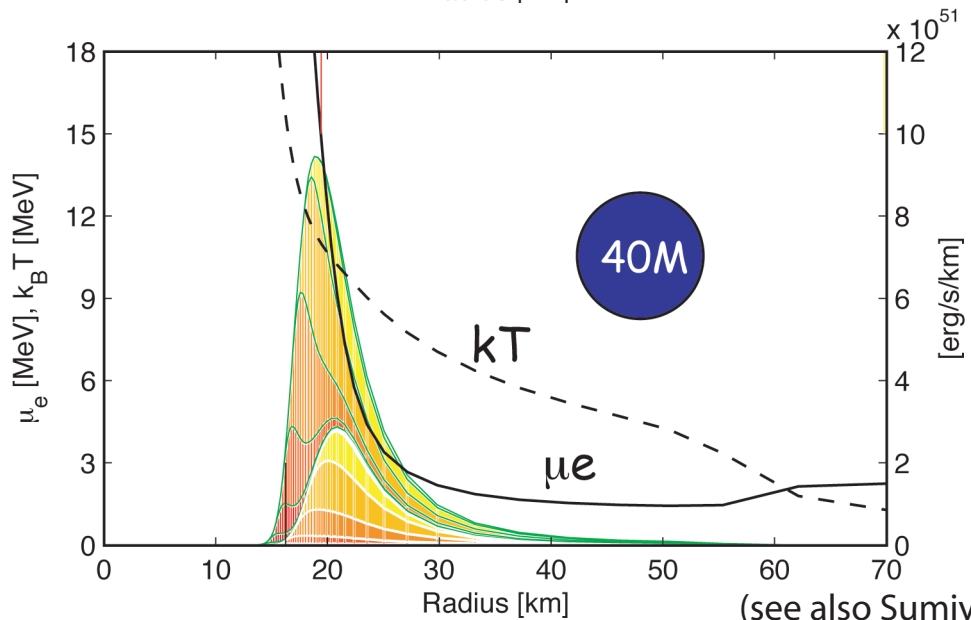
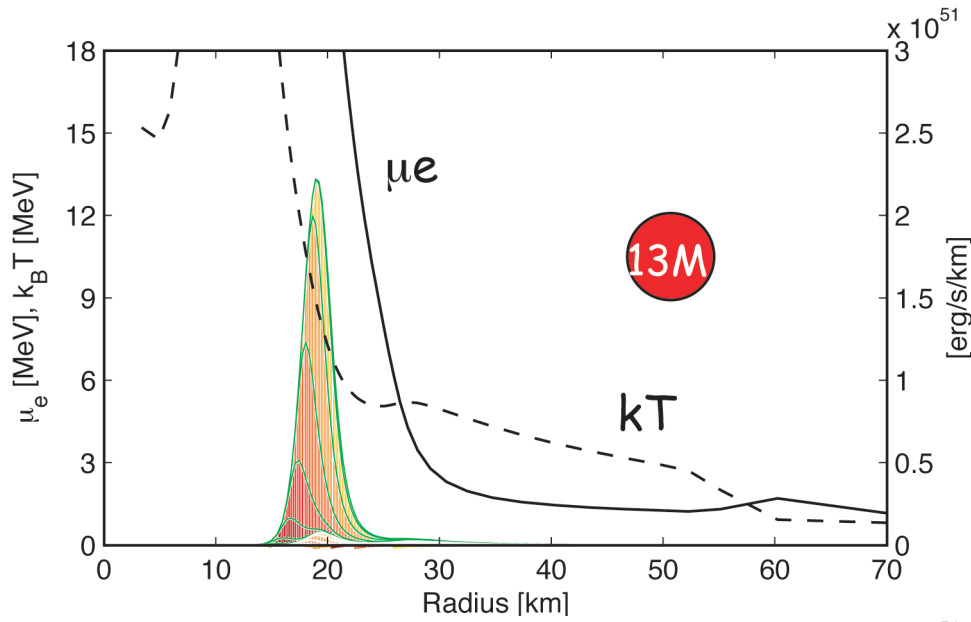
Typical energy hierarchy of neutrino energies:



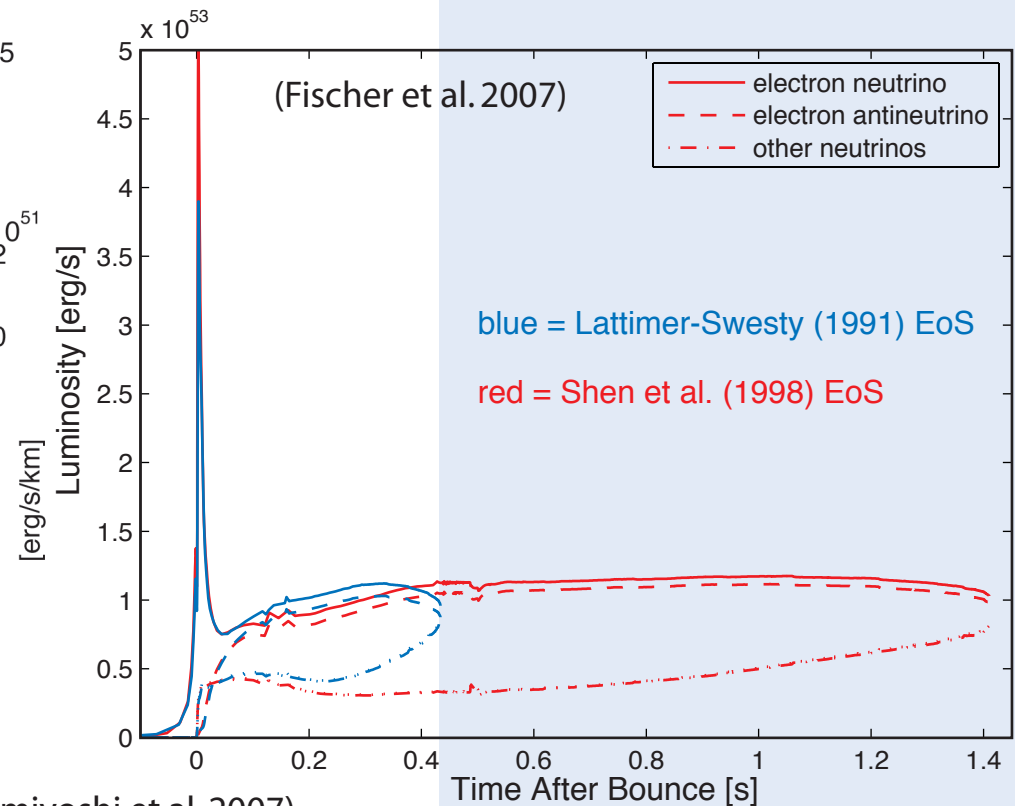
Histogram of electron neutrino emission



PNS evolution & $\nu(\mu/\tau)$ properties

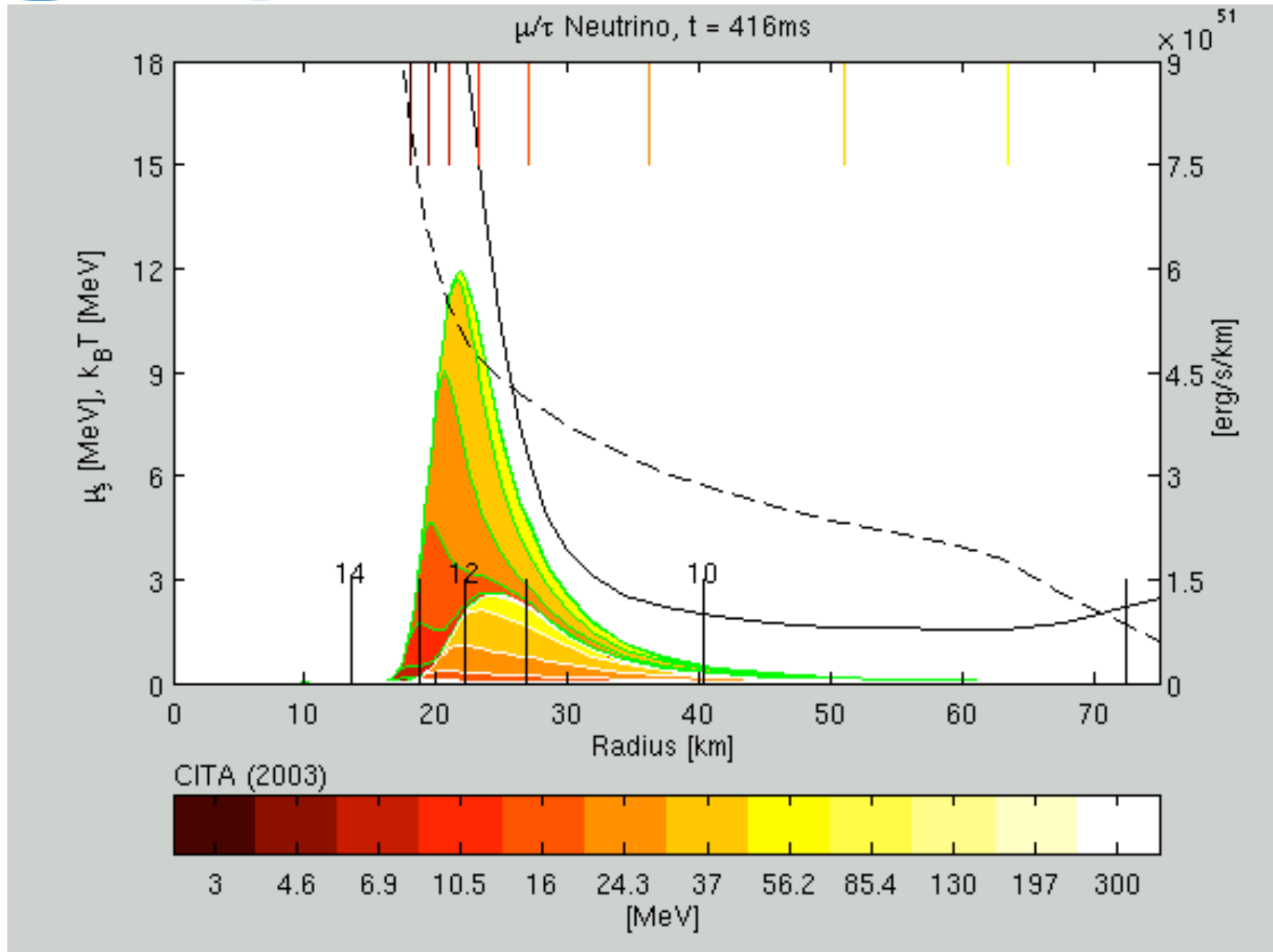


- low mass proto-neutron star (PNS)
--> incompressible accretion
- PNS close to maximum mass
--> hot layers pushed inward

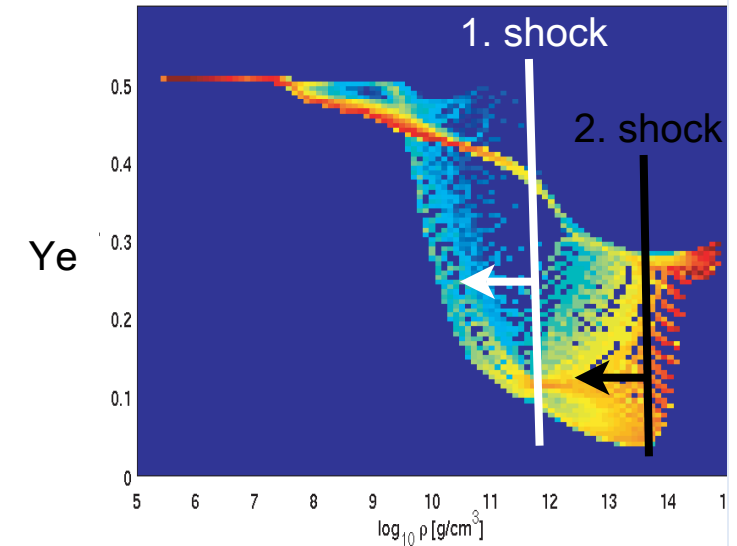
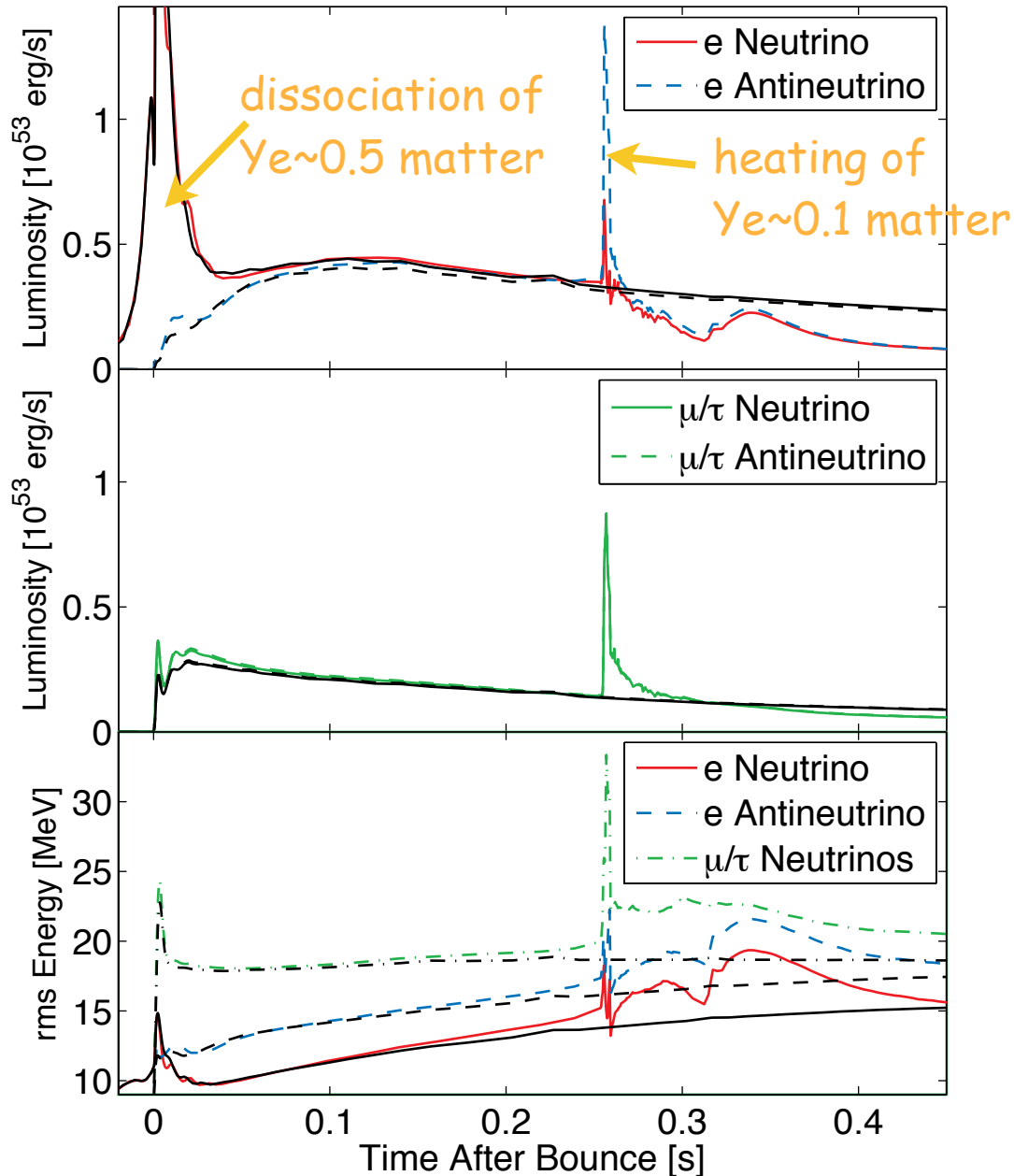


(see also Sumiyoshi et al. 2007)

Histogram of μ/τ neutrino emission



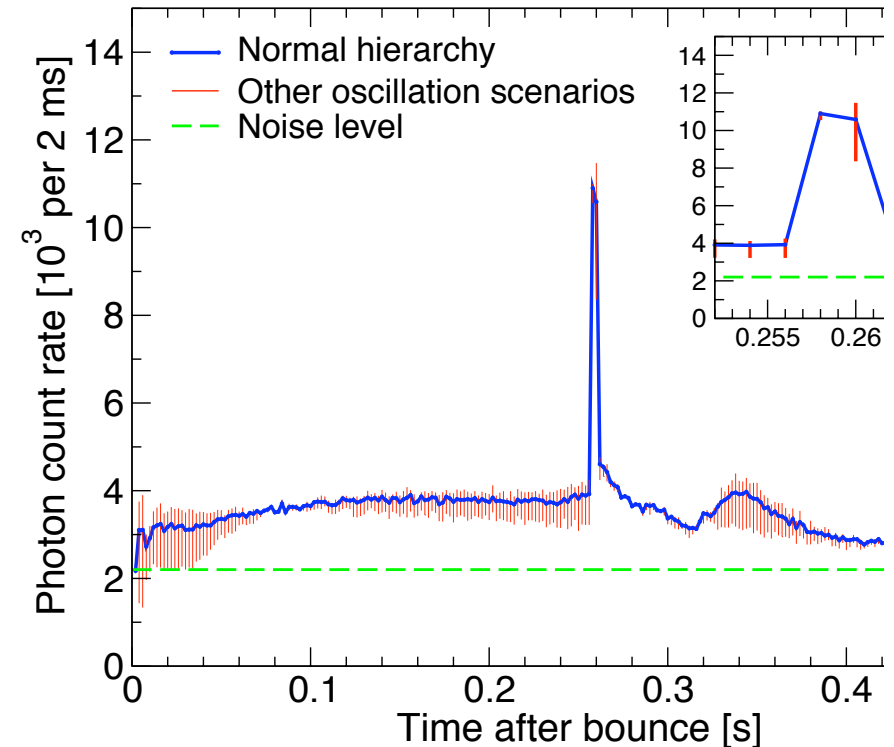
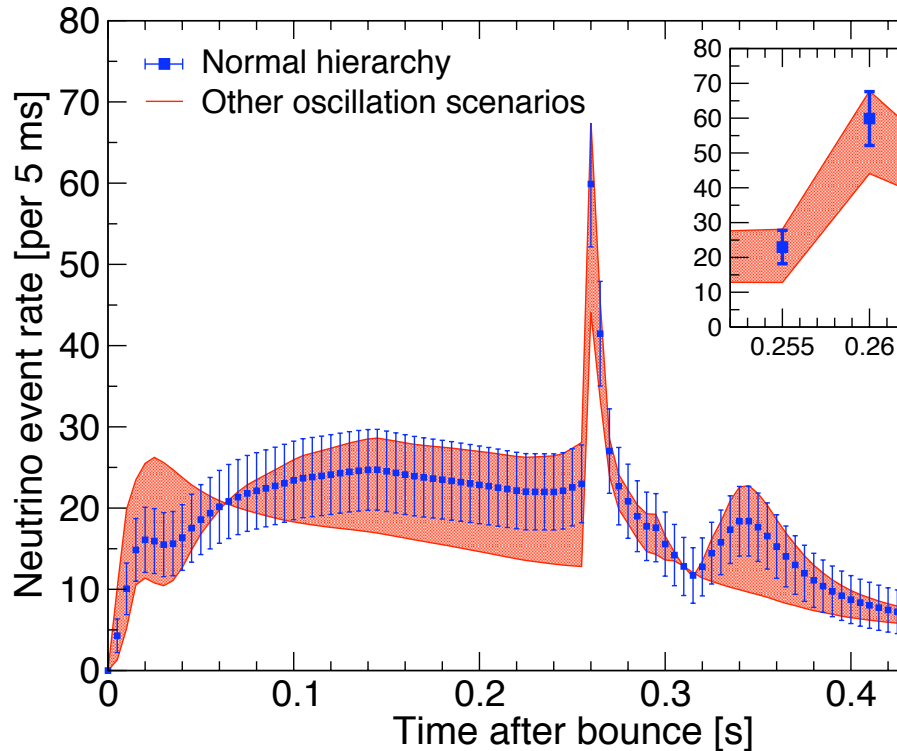
QCD phase transition induced explosions



- Second neutrino peak in all flavours, dominated by anti- ν 's
- Step up in neutrino rms energies

(Sagert, Fischer et al., PRL 2009)

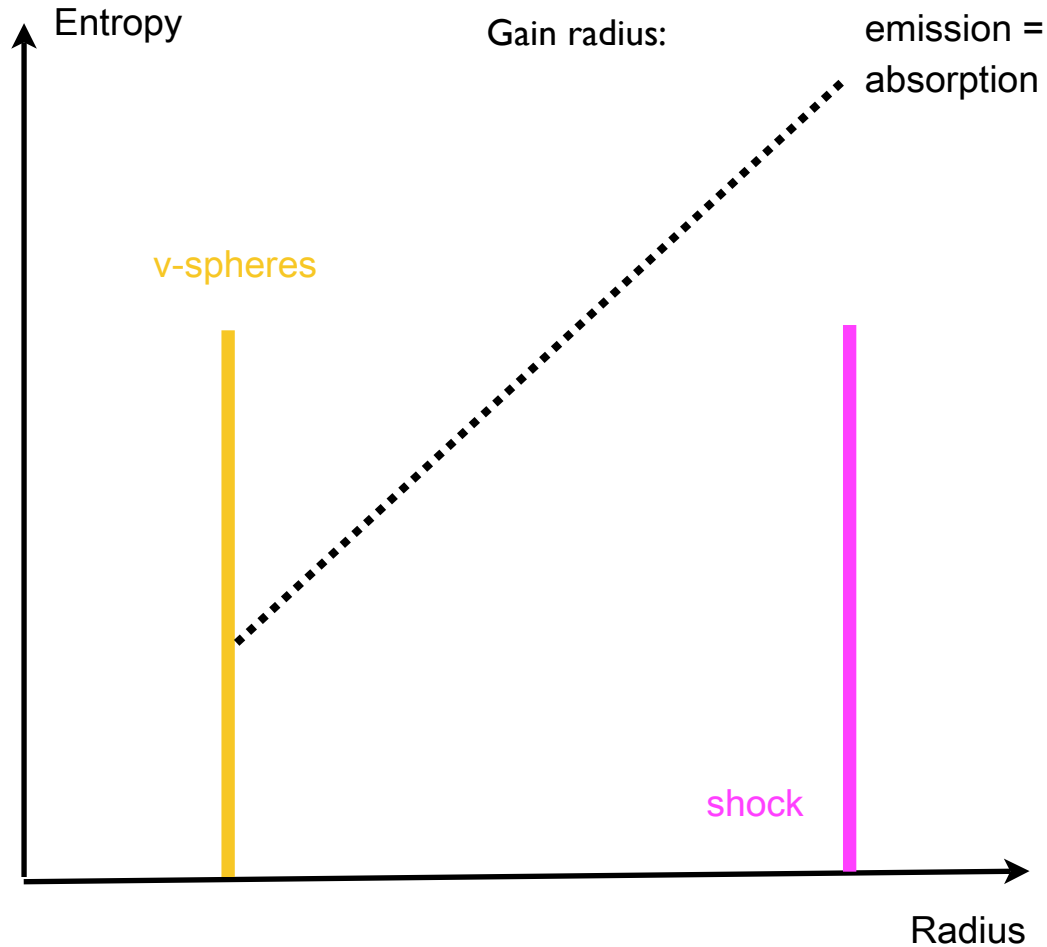
SN as high-density physics laboratory



Neutrino signature from early QCD phase transition
in proto-neutron star (EOS not compatible with 1.97
Ms neutron star! (Demorest et al. 2010))

(Dasgupta et al. 2010)

Physics for the postbounce phase



There are no heavy nuclei between the neutrinosphere and the shock!

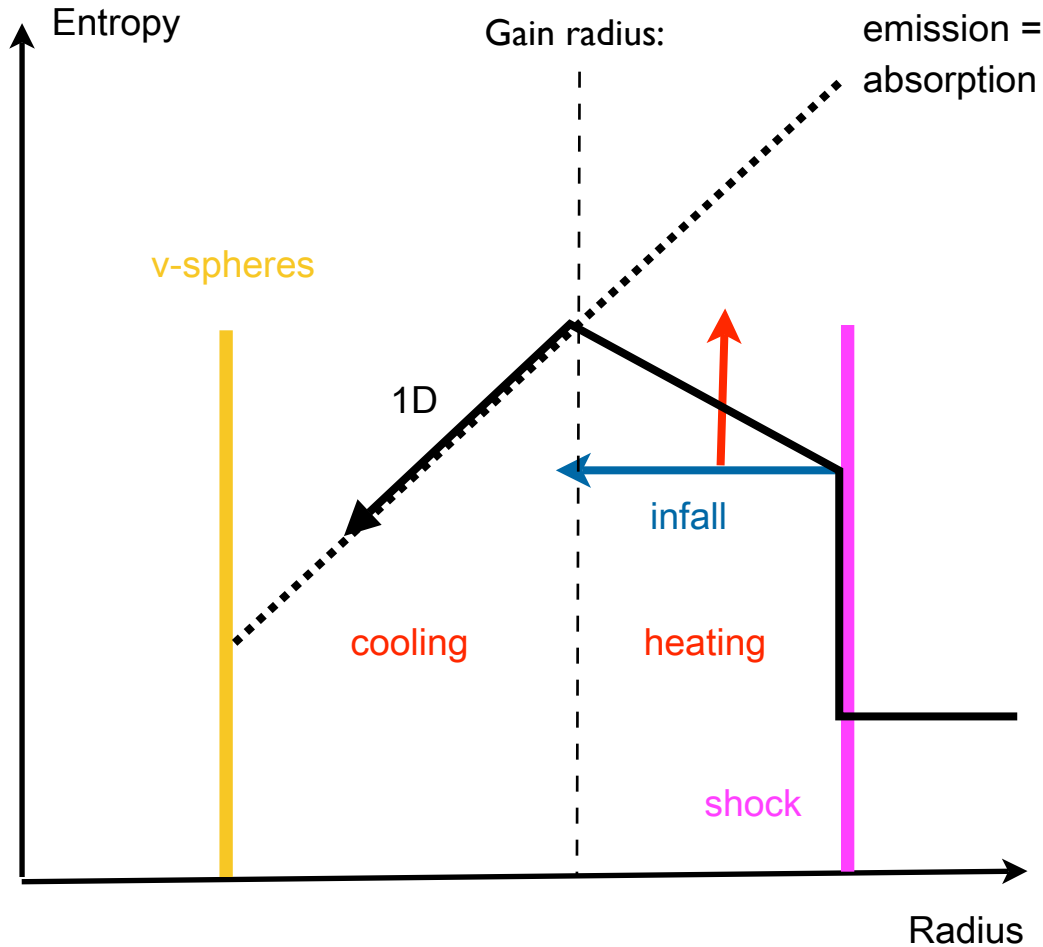
Key input:

- high-density EOS
- fluid instabilities
- neutrino transport, (and oscillations?)
- magnetic fields

Fluid overturn is essential for ν -heating efficiency

(Herant et al. 1994)

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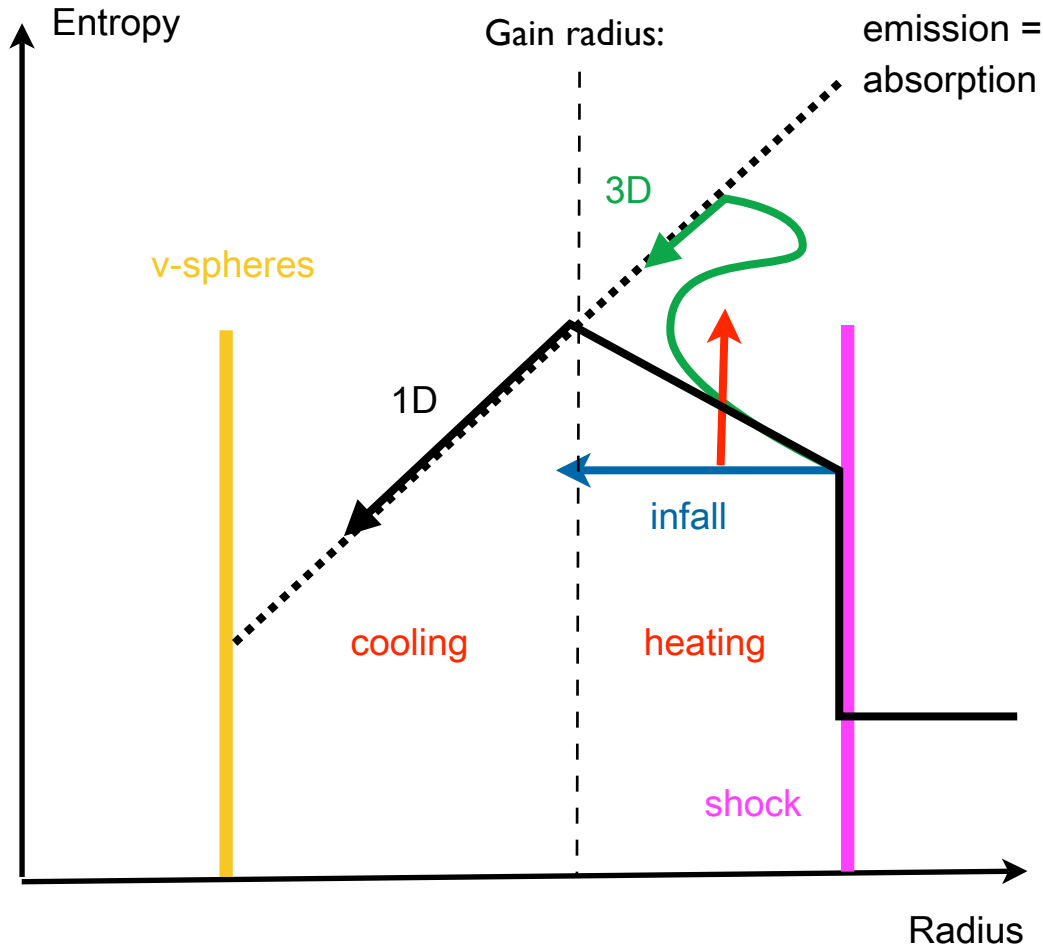
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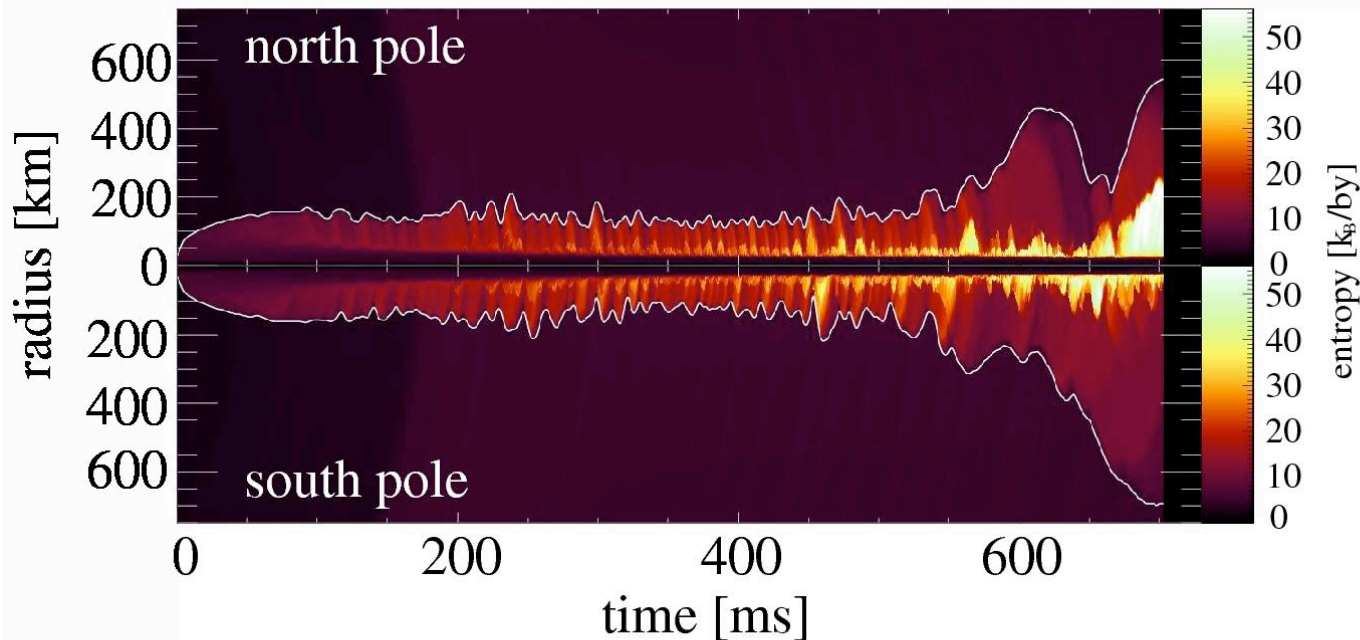
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2D Models in axisymmetry



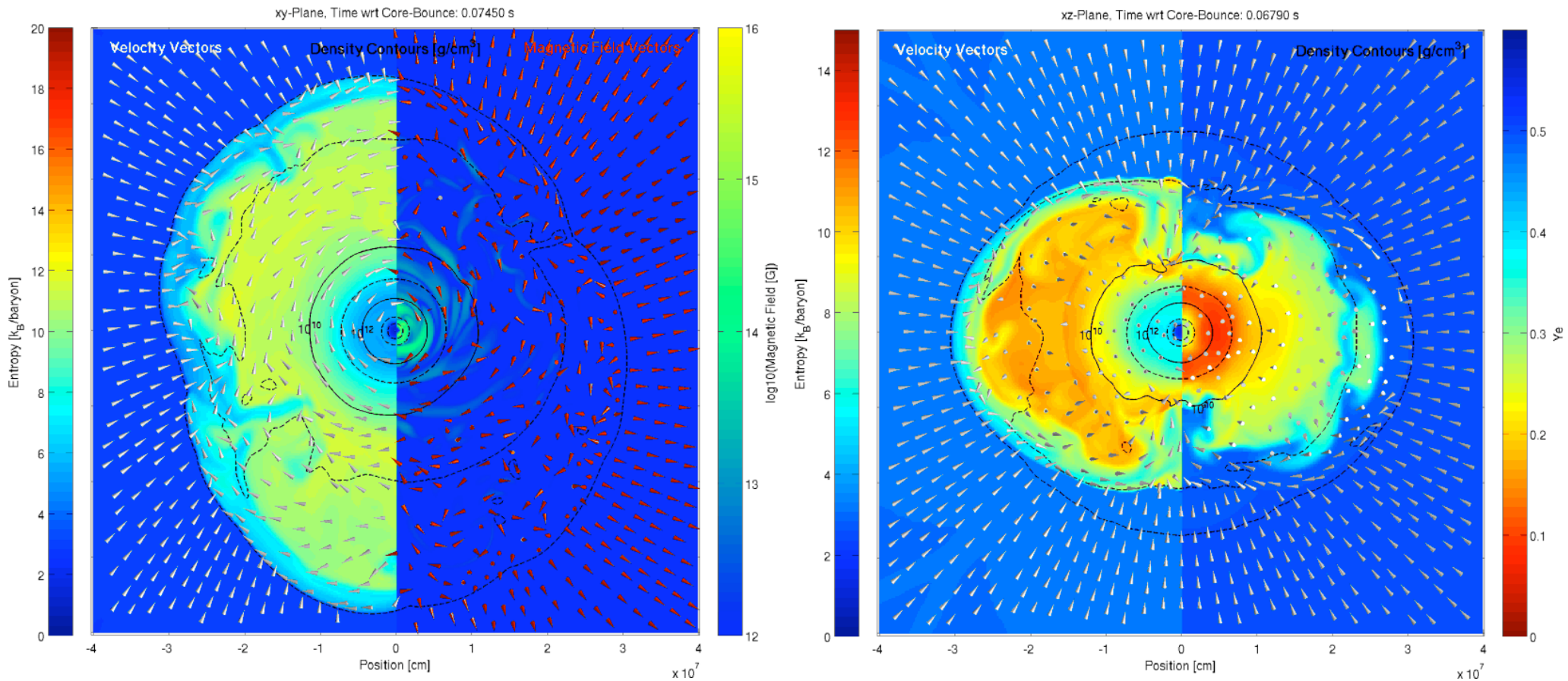
Marek & Janka 2009

- Spectral neutrino transport
- Ray-by-ray approach couples transport in angular wedges to 2D hydrodynamics

Talks by

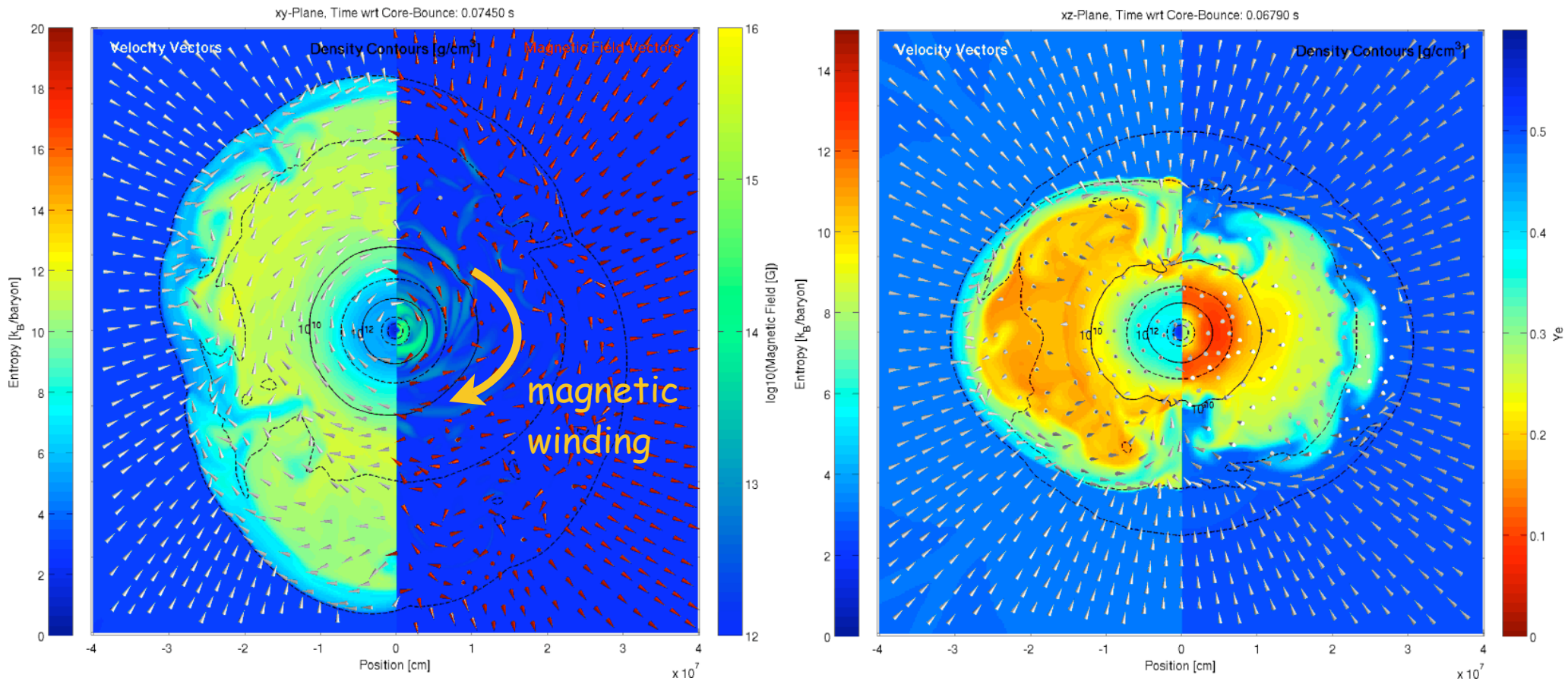
- Steve Bruenn
 - Christian Ott
 - Bernhard Müller
 - Thomas Janka
-
- Neutrino heating
 - Standing Accretion Shock Instability
 - Input physics
 - Asymmetries
 - Grid/Coordinates
 - Relativity / $O(v/c)$

FISH code: 3D magneto-hydrodynamics



- Convective turnover in 2D is restricted to toroidal shapes!
- Tube-shaped downstreams and broad upflows cannot be modelled in 2D.
- Fluid instabilities and coupling to magnetic are intrinsically 3D.

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Solving the Boltzmann equation



$$\begin{aligned}
 & \frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m} + \Gamma \left(\frac{1}{r} - \frac{\partial \alpha}{\alpha \partial r} \right) \frac{\partial [(1 - \mu^2) F]}{\partial \mu} \\
 & + \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) \frac{\partial [\mu (1 - \mu^2) F]}{\partial \mu} \\
 & + \left[\mu^2 \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) - \frac{1u}{r c} - \mu \Gamma \frac{\partial \alpha}{\alpha \partial r} \right] \frac{1}{E^2} \frac{\partial (E^3 F)}{\partial E} \\
 & = \frac{j}{\rho} - \tilde{\chi} F + \frac{1}{h^3 c^4} E^2 \int d\mu' R_{is}(\mu, \mu', E) F(\mu', E) \\
 & - \frac{1}{h^3 c^4} E^2 F \int d\mu' R_{is}(\mu, \mu', E) \\
 & + \frac{1}{h^3 c^4} \left[\frac{1}{\rho} - F(\mu, E) \right] \int E'^2 dE' d\mu' \tilde{R}_{nes}^{in}(\mu, \mu', E, E') F(\mu', E) \\
 & - \frac{1}{h^3 c^4} F(\mu, E) \int E'^2 dE' d\mu' \tilde{R}_{nes}^{out}(\mu, \mu', E, E') \left[\frac{1}{\rho} - F(\mu', E') \right]
 \end{aligned}$$

$$\frac{\partial Y_e}{\partial t} = -\frac{2\pi m_B}{h^3 c^2} \int E^2 dE d\mu \left(\frac{j}{\rho} - \tilde{\chi} F \right) \quad \frac{\partial e}{\partial t} = \dots \quad \frac{\partial u}{\partial t} = \dots$$

Evolution of specific neutrino distr. function:

$$F(t, m, \mu, E) = f(t, r, \mu, E) / \rho$$

=> 3D implicit problem

Comoving metric:

$$\begin{aligned}
 ds^2 &= -\alpha^2 dt^2 + \left(\frac{1}{\Gamma} \frac{\partial r}{\partial a} \right)^2 \\
 &+ r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)
 \end{aligned}$$

Stress-energy tensor:

$$\begin{aligned}
 T^{tt} &= \rho(1 + e + J) \\
 T^{ta} = T^{at} &= \rho H \\
 T^{aa} &= p + \rho K \\
 T^{\vartheta\vartheta} = T^{\varphi\varphi} &= p + \frac{1}{2} \rho (J - K)
 \end{aligned}$$

(Mezzacappa & Bruenn 1993, Liebendörfer 2000, Liebendörfer et al. 2004)

Pitfalls of multi-D Boltzmann ν -transport



Boltzmann transport:

- One fluid element contains
4 ν types \times 20 energies \times 100 angles = 8000 variables
- At a resolution of 1000^3 zones
--> 64TB per time step

Hydrodynamics:

- One fluid element contains ~ 10 variables
- At a resolution of 1000^3 zones
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Pitfalls of multi-D Boltzmann ν -transport

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Compression of Fermi-gas:

$$\frac{dF}{dt} - \frac{1}{3E^2} \frac{\partial}{\partial E} (E^3 \rho F) \frac{d}{dt} \left(\frac{1}{\rho} \right) - \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{c\lambda}{3} \frac{\partial F}{\partial r} \right) = \left(\frac{dF}{dt} \right)_{\text{collision}}$$

de pdV diffusion = interactions

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difficult energy-terms
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Diffusion limit:

$$\frac{\lambda}{3} \frac{\partial F}{\partial r} \ll F, \quad \frac{H}{cJ} \sim 10^{-4}, \quad H = \int_{-1}^{+1} F(\mu) \mu d\mu$$

Hydrodynamics:

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Inaccurate fluxes in
diffusion-regime due to
large cancellations in
angle integral!

There is no perfect transport algorithm...



	Diffusive regime	Semi-transparent	Transparent regime
Boltzmann solver	Truncation errors in flux		Inefficient ang. resol.
Flux-limited diffusion		Flux-factor estimated	Flux-factor unknown
Ray-tracing	Short mean free path	Limited by reaction rates	

The ideal algorithm combines the three green fields! However, it might be too complicated. Alternatives:

Interpolate from diffusive regime to transparent regime:

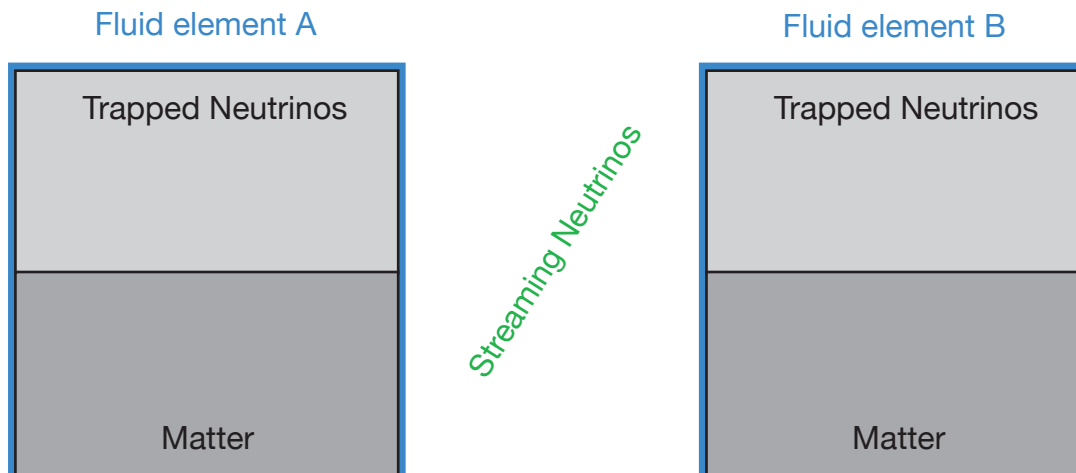
- Multi-Group flux-limited diffusion
Flux factor unknown in transparent regime!
- New: Isotropic diff. source approximation
 $f = f(\text{trapped}) + f(\text{streaming})$, separate evolution Eqs.

Isotropic Diffusion Source Approximation

$$D(f) = j - \chi * f$$

$$f = f(\text{trapped}) + f(\text{streaming}) = f_t + f_s$$

Different approx.
for trapped & streaming
neutrino components!



Isotropic Diffusion Source Approximation

$$D(f) = j - \chi * f$$

$$f = f(\text{trapped}) + f(\text{streaming}) = f_t + f_s$$

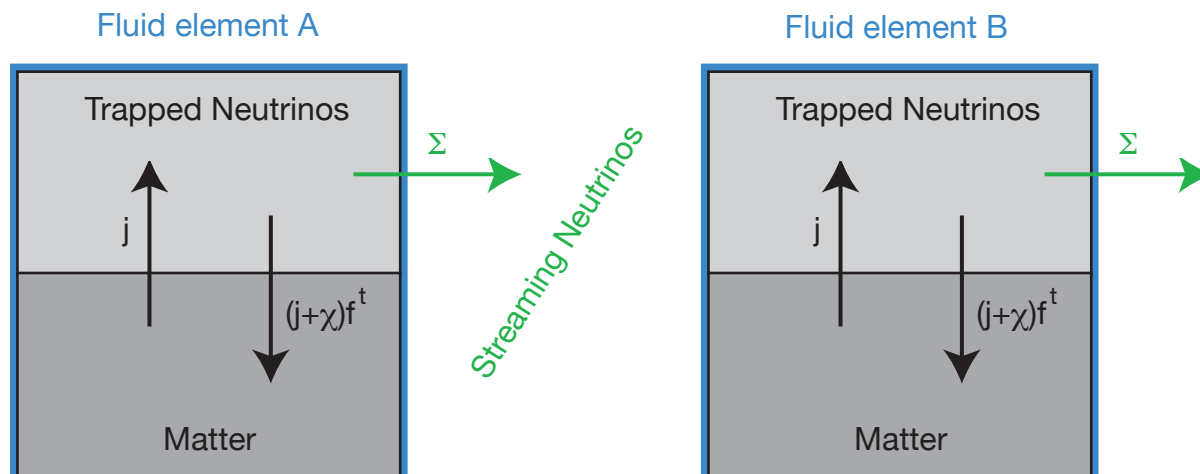
$$D(f_t) = j - \chi * f_t - \Sigma \quad (1)$$

$$D(f_s) = -\chi * f_s + \Sigma \quad (2)$$

Different approx.
for trapped & streaming
neutrino components!

Σ determined by diffusion limit of (1)

- implicit local weak interaction solver
- advection-diffusion problem



Isotropic Diffusion Source Approximation

$$D(f) = j - \chi^* f$$

$$f = f(\text{trapped}) + f(\text{streaming}) = f^t + f^s$$

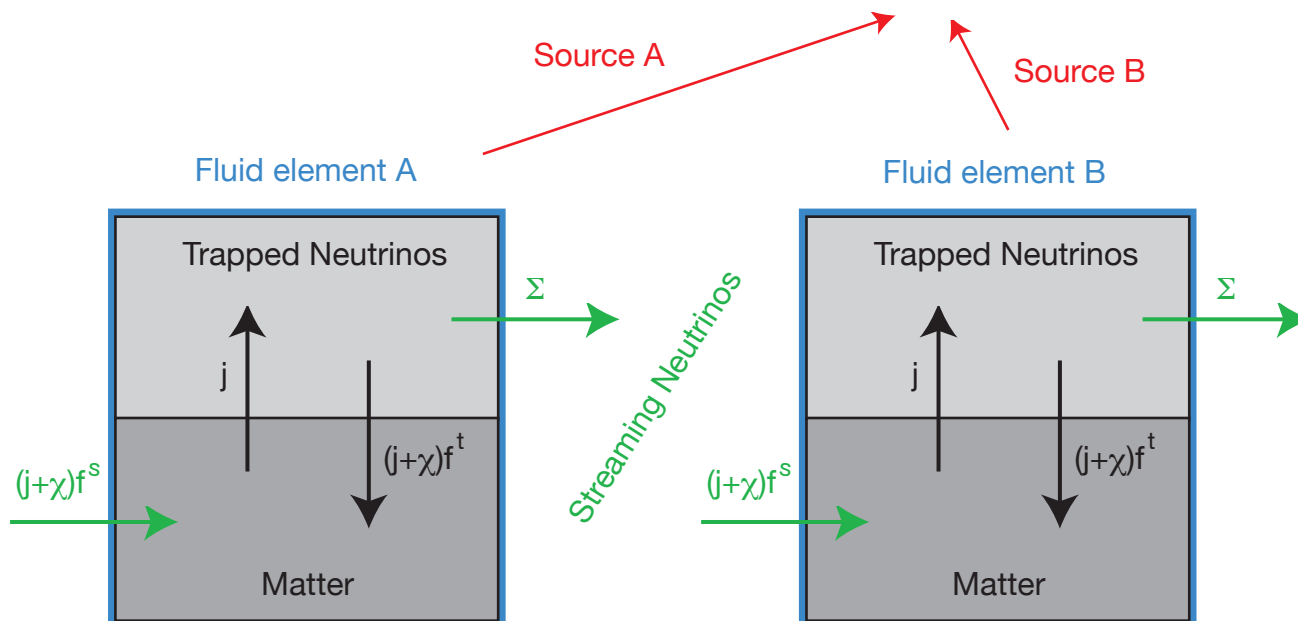
$$D(f^t) = j - \chi^* f^t - \Sigma \quad (1)$$

$$D(f^s) = -\chi^* f^s + \Sigma \quad (2)$$

Different approx.
for trapped & streaming
neutrino components!

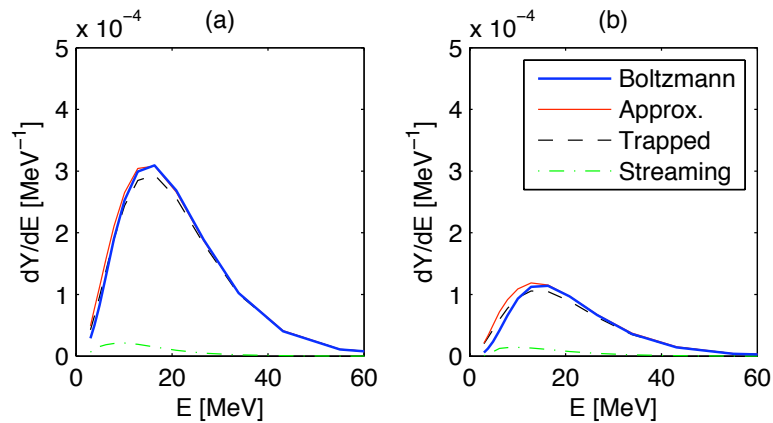
Σ determined by diffusion limit of (1)

Stationary state approx. for (2) --> **Poisson Eq.**



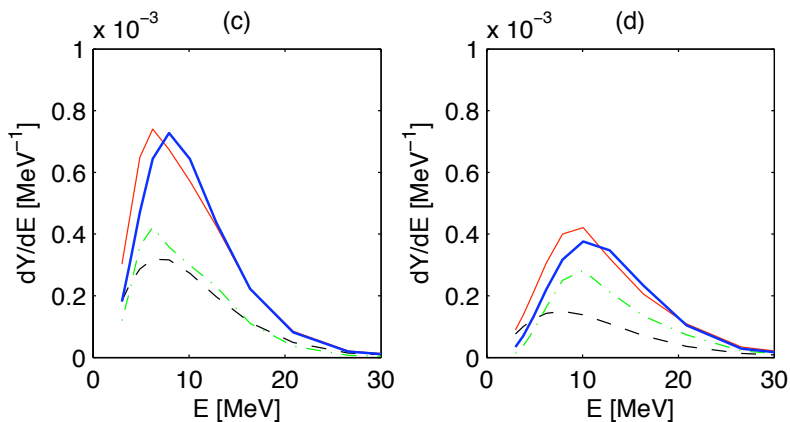
- implicit local weak interaction solver
- advection-diffusion problem
- 20+1 Poisson solves per step
- Geometrical analysis of neutrino-spheres

Comparison of IDSA Spectra



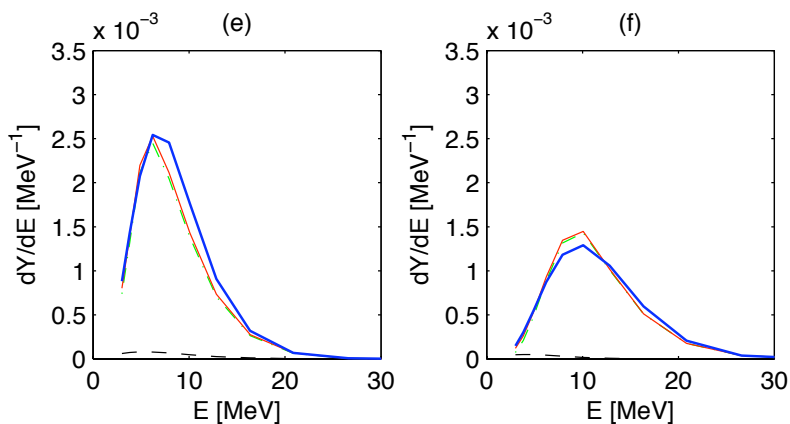
at 40 km radius
(trapped regime)

Trapped neutrinos
dominate spectrum



at 80 km radius
(semi-transparent)

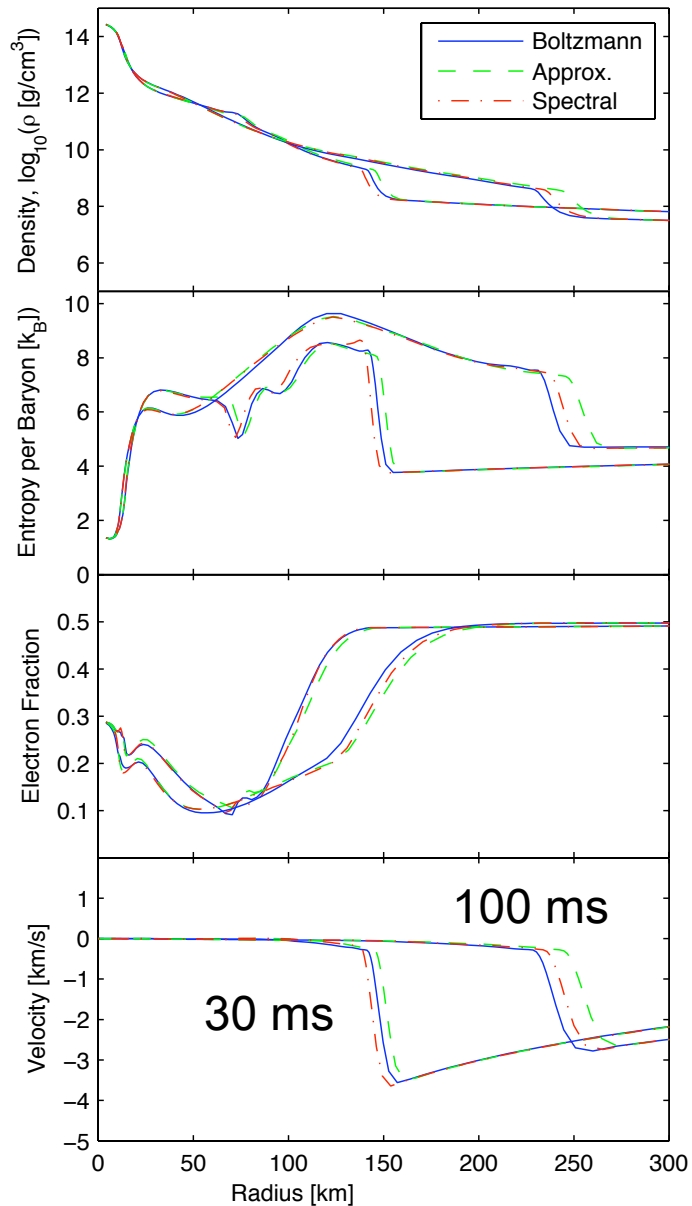
Trapped *and*
streaming neutrinos
form spectrum



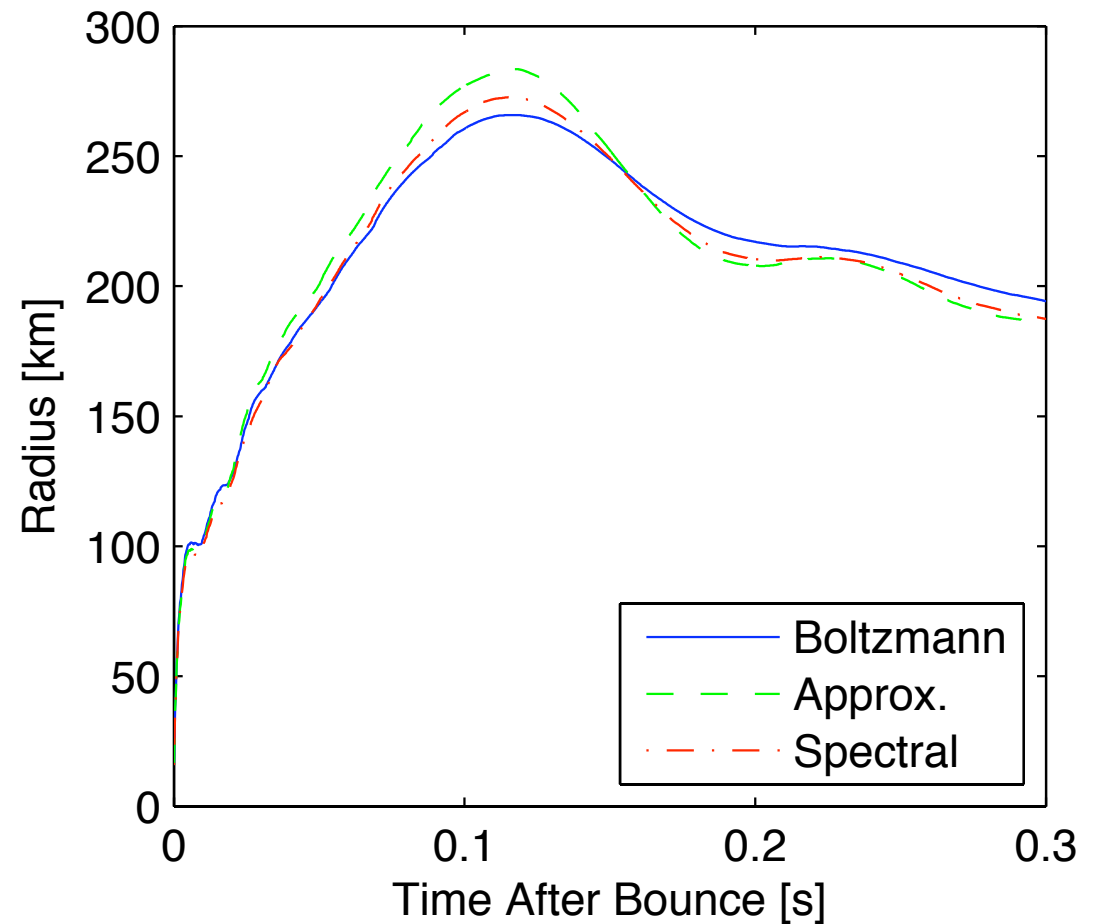
at 160 km radius
(free streaming)

Streaming neutrinos
dominate spectrum

Comparison of Hydrodynamical Evolution



Evolution of shock radius as function of time



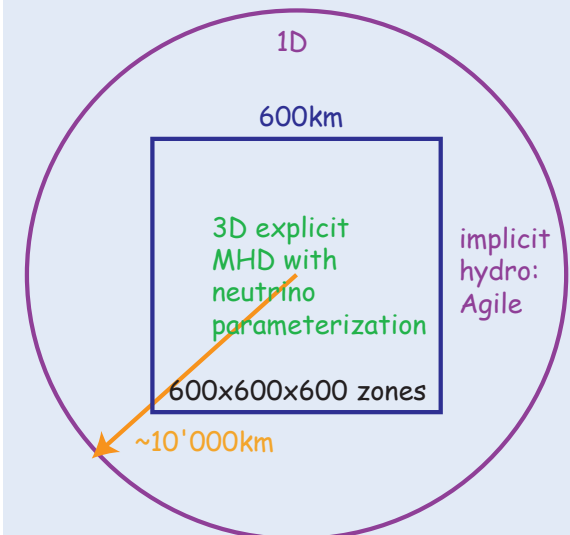
prel. Agile & IDSA available by email
to matthias.liebendoerfer@unibas.ch

From FISH to ELEPHANT...

Elegant parallel hydrodynamics with
approximate neutrino transport

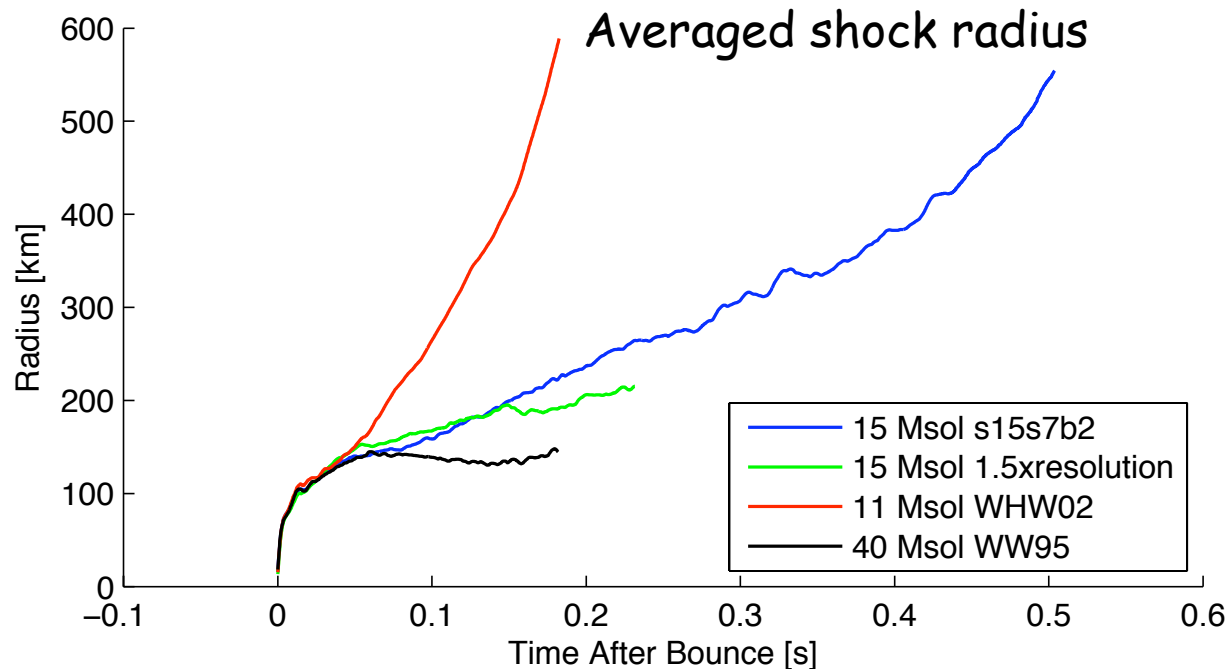


- Lattimer-Swesty EoS
- Effective GR potential
- constrained $\text{div}(B)$
- 2nd order TVD
- IDSA for e-flavour ν 's
- Leakage for μ/τ ν 's



(Whitehouse/Käppeli et al.)

3D supernova model with IDSA



- The 11 solar mass progenitor run shows positive velocities and produces an explosion
- The 15 solar mass progenitor has not (yet) developed positive velocities, even if the shock radius increases.

Runs more optimistic than (Marek & Janka 2009) 2D models.

Runs more pessimistic than (Hix et al. 2010) 2D models.

Convergence test: blue run perhaps still too optimistic?

Entropy

Ye

